

# A Remark On A Nonlinear System Arising In A Theory Of Thermal Explosion\*

Huda Ibrahim Alsayed<sup>†</sup>, Elaf Ali Alsisi<sup>‡</sup>, Divya Adalja<sup>§</sup>, Ramappa Manjunatha<sup>¶</sup>  
Udandarao Sarat<sup>||</sup>, Damanjeet Aulakh<sup>\*\*</sup>, Kamaljit Kaur<sup>††</sup>, Shivakrishna Dasi<sup>‡‡</sup>  
Hessan Mohammad<sup>§§¶</sup>

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## Abstract

The purpose of this paper is to study the existence and multiplicity of positive solutions for a mathematical model of thermal explosion which is described by the system

$$\begin{cases} -\Delta u = \lambda f(v), & x \in \Omega, \\ -\Delta v = \lambda g(u), & x \in \Omega, \\ \mathbf{n} \cdot \nabla u + a(u)u = 0, & x \in \partial\Omega, \\ \mathbf{n} \cdot \nabla v + b(v)v = 0, & x \in \partial\Omega, \end{cases}$$

where  $\Omega$  is a bounded smooth domain of  $\mathbb{R}^N$ ,  $\Delta$  is the Laplacian operator,  $\lambda > 0$  is a parameter,  $f, g$  are  $C^1$  strictly increasing functions that have a combined sublinear effect at  $\infty$ , and  $a, b : [0, \infty) \rightarrow (0, \infty)$  are strictly increasing  $C^1$  functions. We establish our existence and multiplicity results by the method of sub- and supersolutions.

## 1 Introduction

A classical problem in combustion theory is a model of thermal explosion which occurs due to a spontaneous ignition in a rapid combustion process. In this paper, we consider a model involving a nonlinear boundary heat loss which is not a very typical one in classical combustion theory, but is relevant to some more applications (see [4, 10, 12, 5] for details). The model reads as:

$$\begin{cases} \theta(t) - \Delta\theta = \lambda f(\eta), & (t, x) \in (0, \infty) \times \Omega, \\ \eta(t) - \Delta\eta = \lambda g(\theta), & (t, x) \in (0, \infty) \times \Omega, \\ \mathbf{n} \cdot \nabla\theta + a(\theta)\theta = 0, & (t, x) \in (0, \infty) \times \partial\Omega, \\ \mathbf{n} \cdot \nabla\eta + b(\eta)\eta = 0, & (t, x) \in (0, \infty) \times \partial\Omega, \\ \theta(0, x) = 0 = \eta(0, x). \end{cases} \quad (1)$$

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<sup>†</sup>Accounting Department, Faculty of Business School, King Khalid University, P.O. Box 9004, Abha 61413, Saudi Arabia

<sup>‡</sup>Business Informatics Department, Business College, King Khalid University, Abha, Saudi Arabia

<sup>§</sup>Department of Sciences, Vivekananda Global University, Jaipur, Rajasthan-303012, India

<sup>¶</sup>Department of Data analytics and Mathematical Sciences, School of Sciences, JAIN (Deemed to be University), Bangalore, Karnataka

<sup>||</sup>Department of Mathematics Raghu Engineering College, Dakamarri-531162, Andhra Pradesh, India

<sup>\*\*</sup>Centre for Research Impact & Outcome, Chitkara University Institute of Engineering and Technology Chitkara University, Rajpura, 140401, Punjab, India

<sup>††</sup>Department of Applied Sciences, Chandigarh Engineering College, Chandigarh Group of Colleges-Jhanjeri, Mohali, Pin 140307, Punjab, India

<sup>‡‡</sup>NIMS School of Computing Science and Artificial Intelligence, NIMS University Rajasthan, Jaipur, India

<sup>§§</sup>Department of computers Techniques engineering, College of technical engineering, The Islamic University, Najaf, Iraq

<sup>¶¶</sup>Department of computers Techniques engineering, College of technical engineering, The Islamic University of Al Diwaniyah, Al Diwaniyah, Iraq

Here  $\theta, \eta$  are the appropriately scaled temperature in a bounded smooth domain  $\Omega \subset \mathbb{R}^N$ ,  $N \geq 1$ , and  $f, g$  are the normalized reaction rate. We assume that  $f, g$  satisfy the following assumptions:

(H1)  $f, g \in C([0, \infty))$  are strictly increasing functions,

(H2)  $\lim_{s \rightarrow \infty} \frac{f(Ag(s))}{s} = 0$ , for all  $A > 0$ .

On the  $C^2$  boundary  $\partial\Omega$ , with the outward unit normal denoted by  $\mathbf{n}$ , the heat-loss parameters  $a(\theta), b(\eta)$  are assumed to satisfy the following hypothesis:

(H3)  $a, b : [0, \infty) \rightarrow (0, \infty)$  are strictly increasing bounded  $C^1$  functions.

Physically this assumption means that a heat loss through the boundary always exists and increases linearly with the temperature even in the small temperature regime.

A bifurcation (or scaling) parameter  $\lambda > 0$  can be associated with the size of domain  $\Omega$  in (1) which grows linearly as the measure of  $\Omega$  increases. It is well known that, after normalizing for the size of  $\Omega$ , the long term behavior of solution of the system (1) is close to the solution of the time-independent system:

$$\begin{cases} -\Delta u = \lambda f(v), & x \in \Omega, \\ -\Delta v = \lambda g(u), & x \in \Omega, \\ \mathbf{n} \cdot \nabla u + a(u)u = 0, & x \in \partial\Omega, \\ \mathbf{n} \cdot \nabla v + b(v)v = 0, & x \in \partial\Omega. \end{cases} \quad (2)$$

The motivation for this study comes from the work in [7] where the authors established the existence, uniqueness and multiplicity of positive solutions for certain range of  $\lambda$  for the single equation of the form

$$\begin{cases} -\Delta u = \lambda f(u), & x \in \Omega, \\ \mathbf{n} \cdot \nabla u + a(u)u = 0, & x \in \partial\Omega. \end{cases}$$

Here we extend this study to Laplacian system of the form (2). In [1], Ali-Shivaji-Ramaswamy discussed the existence of multiple positive solutions to such systems with Dirichlet boundary conditions. One can refer to [3, 8] for some recent existence and uniqueness results of elliptic problems with nonlinear boundary conditions.

## 2 Existence Results

In this section, we shall establish our existence results via the method of sub - supersolution. A pair of nonnegative functions  $(\psi_1, \psi_2) \in W^{1,2} \cap C(\bar{\Omega}) \times W^{1,2} \cap C(\bar{\Omega})$  and a pair  $(z_1, z_2) \in W^{1,2} \cap C(\bar{\Omega}) \times W^{1,2} \cap C(\bar{\Omega})$  are called a subsolution and supersolution of the system (2) if they satisfy

$$\begin{cases} -\Delta \psi_1 \leq \lambda f(\psi_2), & x \in \Omega, \\ -\Delta \psi_2 \leq \lambda g(\psi_1), & x \in \Omega, \\ \mathbf{n} \cdot \nabla \psi_1 + a(\psi_1)\psi_1 \leq 0, & x \in \partial\Omega, \\ \mathbf{n} \cdot \nabla \psi_2 + b(\psi_2)\psi_2 \leq 0, & x \in \partial\Omega, \end{cases} \quad (3)$$

and

$$\begin{cases} -\Delta z_1 \geq \lambda f(z_2), & x \in \Omega, \\ -\Delta z_2 \geq \lambda g(z_1), & x \in \Omega, \\ \mathbf{n} \cdot \nabla z_1 + a(z_1)z_1 \geq 0, & x \in \partial\Omega, \\ \mathbf{n} \cdot \nabla z_2 + b(z_2)z_2 \geq 0, & x \in \partial\Omega, \end{cases} \quad (4)$$

respectively. It is well known that if there exist sub and supersolutions  $(\psi_1, \psi_2)$  and  $(z_1, z_2)$  respectively of the system (2) such that  $(\psi_1, \psi_2) \leq (z_1, z_2)$ . Then the system (2) has a solution  $(u, v)$  such that  $(u, v) \in [(\psi_1, \psi_2), (z_1, z_2)]$  ( see [2, 6] ).

By strict sub and super-solutions we understand functions  $(\psi_1, \psi_2)$  and  $(z_1, z_2)$  for which strict inequalities (3) and (4) hold.

Our multiplicity results are obtained by constructing sub and super-solution pairs that satisfy the following lemma.

**Lemma 1** ([6, 9, 11]) *Suppose that the system (2) has a sub-solution  $(\psi_1, \psi_2)$ , a strict super-solution  $(\zeta_1, \zeta_2)$ , a strict sub-solution  $(w_1, w_2)$ , and a super-solution  $(z_1, z_2)$  for the system (2) such that*

$$\begin{aligned} (\psi_1, \psi_2) &\leq (\zeta_1, \zeta_2) \leq (z_1, z_2), \\ (\psi_1, \psi_2) &\leq (w_1, w_2) \leq (z_1, z_2), \end{aligned}$$

and  $(w_1, w_2) \not\leq (\zeta_1, \zeta_2)$ . Then the system (2) has at least three distinct solutions  $(u_i, v_i)$ ,  $i = 1, 2, 3$  such that

$$(u_1, v_1) \in [(\psi_1, \psi_2), (\zeta_1, \zeta_2)], \quad (u_2, v_2) \in [(w_1, w_2), (z_1, z_2)]$$

and

$$(u_3, v_3) \in [(\psi_1, \psi_2), (z_1, z_2)] \setminus \left( [(\psi_1, \psi_2), (\zeta_1, \zeta_2)] \cup [(w_1, w_2), (z_1, z_2)] \right).$$

To precisely state our existence result we consider the unique classical solution  $e_r$  of the following linear elliptic problem

$$\begin{cases} -\Delta e_r = 1, & x \in \Omega, \\ \mathbf{n} \cdot \nabla e_r + r_0 e_r = 0, & x \in \partial\Omega, \end{cases}$$

for  $r = a, b$ , where  $r_0 = r(0)$ . Then we establish the following theorem.

**Theorem 1** *Let (H1)–(H3) hold and  $f(0)$  or  $g(0)$  be strictly positive. Then the system (2) has a positive solution  $(u, v)$  for all  $\lambda > 0$ .*

**Proof.** It is easy to see that  $(\psi_1, \psi_2) = (0, 0)$  is a subsolution of the system (2). We now construct the supersolution  $(z_1, z_2)$ . Let  $(z_1, z_2) = (C_\lambda e_a, \lambda g(C_\lambda \|e_b\|_\infty) e_b)$ , where  $C_\lambda$  is a large number to be chosen later. We shall verify that  $(z_1, z_2)$  is a supersolution of the system (2) for all  $\lambda > 0$ . By (H2) we can choose  $C_\lambda$  large enough so that

$$C_\lambda \geq \lambda f\left(\lambda g(C_\lambda \|e_b\|_\infty) \|e_b\|_\infty\right),$$

and therefore

$$\begin{aligned} -\Delta z_1 &= C_\lambda \geq \lambda f\left(\lambda g(C_\lambda \|e_b\|_\infty) \|e_b\|_\infty\right) \\ &\geq \lambda f\left(\lambda g(C_\lambda \|e_b\|_\infty) e_b\right) \\ &= \lambda f(z_2) \text{ in } \Omega, \end{aligned}$$

and

$$\begin{aligned} \mathbf{n} \cdot \nabla z_1 + a(z_1) z_1 &\geq C_\lambda \mathbf{n} \cdot \nabla e_a + C_\lambda e_a a_0 \\ &= C_\lambda (\mathbf{n} \cdot \nabla e_a + e_a a_0) \\ &= 0 \text{ on } \partial\Omega. \end{aligned}$$

Next,

$$-\Delta z_2 = \lambda g\left(C_\lambda \|e_b\|_\infty\right) \geq \lambda g\left(C_\lambda e_b\right) = \lambda g(z_1), \text{ in } \Omega,$$

and

$$\begin{aligned} \mathbf{n} \cdot \nabla z_2 + b(z_2) z_2 &\geq \lambda g\left(C_\lambda \|e_b\|_\infty\right) \mathbf{n} \cdot \nabla e_b + \lambda g\left(C_\lambda \|e_b\|_\infty\right) e_b b_0 \\ &= \lambda g\left(C_\lambda \|e_b\|_\infty\right) (\mathbf{n} \cdot \nabla e_b + b_0 e_b) = 0, \text{ on } \partial\Omega, \end{aligned}$$

which implies that  $(z_1, z_2)$  is indeed a positive supersolution of the system (2). Therefore the system (2) has a positive solution for all  $\lambda > 0$ . ■

Our second result concerns with multiplicity of solution for the system (2) and gives an estimate on the parameter  $\lambda$  when such a situation occurs. For positive constants  $a_i, b_i; i = 1, 2$ , define

$$Q_1(a_1, b_1) = \min\left\{\frac{a_1}{f(b_1)}, \frac{b_1}{g(a_1)}\right\}$$

and

$$Q_2(a_2, b_2) = \max\left\{\frac{a_2}{f(b_2)}, \frac{b_2}{g(a_2)}\right\}.$$

Then we establish:

**Theorem 2** Assume  $f(0)$  or  $g(0)$  be positive. Let  $B_R$  be the largest ball of radius  $R$  inscribed in  $\Omega$ , for  $0 < \epsilon < R$ , we define

$$C_1(\Omega) = \inf_{\epsilon} \frac{N}{\epsilon^N} \frac{R^{N-1}}{R - \epsilon},$$

and  $C(\Omega) = C_1(\Omega)\|e_r\|_{\infty}$ , for  $r = a, b$ . Let (H1)–(H3) hold and  $\frac{Q_1}{Q_2} > C(\Omega)$  for some  $a_i, b_i, i = 1, 2$ . Then the system (2) has at least three positive solutions for  $\lambda \in (\lambda_*, \lambda^*)$ , where  $\lambda_* = CQ_2$  and  $\lambda^* = \frac{Q_1}{\|e_r\|_{\infty}}$ , for  $r = a, b$ .

**Proof.** We will establish a pair of subsolutions  $(\psi_1, \psi_2), (w_1, w_2)$  and a pair of supersolutions  $(\zeta_1, \zeta_2), (z_1, z_2)$ , satisfying Lemma 1. Clearly  $(\psi_1, \psi_2) = (0, 0)$  is a subsolution of the system (2).

We next construct a positive supersolution  $(\zeta_1, \zeta_2)$ , of the system (2) when  $\lambda < \frac{Q_1}{\|e_r\|_{\infty}}$ , for  $r = a, b$ . Since  $\lambda < \frac{a_1}{f(b_1)\|e_a\|_{\infty}}$ , we can choose  $\epsilon > 0$  so small that  $\lambda f(b_1) < \frac{a_1}{\epsilon + \|e_a\|_{\infty}}$ . Let

$$(\zeta_1, \zeta_2) = \left(a_1 \frac{e_a + \epsilon}{\|e_a\|_{\infty} + \epsilon}, b_1 \frac{e_b + \epsilon}{\|e_b\|_{\infty} + \epsilon}\right).$$

Then, we have

$$\begin{aligned} -\Delta \zeta_1 &= \frac{a_1}{\epsilon + \|e_a\|_{\infty}} > \lambda f(b_1) \\ &\geq \lambda f\left(b_1 \frac{e_b + \epsilon}{\|e_b\|_{\infty} + \epsilon}\right) \\ &= \lambda f(\zeta_2) \text{ in } \Omega, \end{aligned}$$

and

$$\begin{aligned} \mathbf{n} \cdot \nabla \zeta_1 + a(\zeta_1)\zeta_1 &\geq \frac{a_1}{\epsilon + \|e_a\|_{\infty}} \left(\mathbf{n} \cdot \nabla e_a + (e_a + \epsilon)a_0\right) \\ &= \frac{a_1}{\epsilon + \|e_a\|_{\infty}} (\mathbf{n} \cdot \nabla e_a + a_0 e_a + a_0 \epsilon) \\ &= \frac{a_1 a_0 \epsilon}{\epsilon + \|e_a\|_{\infty}} \\ &> 0 \text{ on } \partial\Omega. \end{aligned}$$

Similar argument shows that  $\zeta_2$  satisfies  $-\Delta \zeta_2 > \lambda g(\zeta_1)$  in  $\Omega$ , and  $\mathbf{n} \cdot \nabla \zeta_2 + b(\zeta_2)\zeta_2 > 0$ .

Next let us construct a strict sub-solution  $(w_1, w_2)$  of the system (2). First note that a system

$$\begin{cases} -\Delta u_D = \lambda f(v_D), & x \in \Omega, \\ -\Delta v_D = \lambda g(u_D), & x \in \Omega, \\ u_D = 0 = v_D, & x \in \partial\Omega, \end{cases}$$

admits a strict sub-solution  $(w_{1D}, w_{2D})$  with  $\|w_{1D}\|_\infty \geq a_2$  and  $\|w_{2D}\|_\infty \geq b_2$  provided  $\lambda < \lambda^*$  (see [1]). Then we have  $(w_1, w_2) \not\leq (\zeta_1, \zeta_2)$ . From Hopf's lemma,  $\mathbf{n} \cdot \nabla w_{iD} < 0$  for  $i = 1, 2$ . Therefore, setting  $w_1 = w_{1D}$  and  $w_2 = w_{2D}$  we obtain a strict sub-solution for the system (2). for  $\lambda > \lambda_*$ .

Let  $(z_1, z_2)$  be the super solution as in the proof of Theorem 2 Further,  $w_i, \zeta_i \leq z_i, i = 1, 2$  for  $C_\lambda$  large. Hence there exist positive solutions  $(u_i, v_i), i = 1, 2, 3$  such that

$$(u_1, v_1) \in [(\psi_1, \psi_2), (\zeta_1, \zeta_2)], (u_2, v_2) \in [(w_1, w_2), (z_1, z_2)]$$

and

$$(u_3, v_3) \in [(\psi_1, \psi_2), (z_1, z_2)] \setminus \left( [(\psi_1, \psi_2), (\zeta_1, \zeta_2)] \cup [(w_1, w_2), (z_1, z_2)] \right).$$

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## References

- [1] J. Ali, R. Shivaji and M. Ramaswamy, Multiple positive solutions for classes of elliptic systems with combined nonlinear effects, *Diff. and Int. Eqns.*, 19(2006), 669–680.
- [2] H. Amann, Fixed point equations and nonlinear eigenvalue problems in ordered Banach spaces, *SIAM Rev.* 18(1976), 620–709.
- [3] D. Butler, E. Ko, E.K. Lee and R. Shivaji, Positive radial solutions for elliptic equations on exterior domains with nonlinear boundary conditions, *Commu. Pure. Appl. Anal*, 13(2014), 2713–2731.
- [4] D. A. Frank-Kamenetskii, *Diffusion and Heat Transfer in Chemical Kinetics*, Plenum Press, New York, 1969.
- [5] P. Gordon, E. Ko and R. Shivaji, Multiplicity and uniqueness of positive solutions for elliptic equations with nonlinear boundary conditions arising in a theory of thermal explosion, *Nonlinear Anal. Real World Appl.*, 15(2014), 51–57.
- [6] F. Inkmann, Existence and multiplicity theorems for semilinear elliptic equations with nonlinear boundary conditions, *Indiana Univ. Math. J.*, 31(1982), 213–221.
- [7] E. Ko and S. Prashanth, Positive solutions for elliptic equations in two dimensionals arising in a theory of thermal explosion, *Taiwanese. J. Math*, 19(2015), 1759–1775.
- [8] E. Ko, M. Ramaswamy and R. Shivaji, Uniqueness of positive radial solutions for a class of semipositone problems on the exterior domain of a ball, *J. Math. Anal, Appl*, 423(2015), 399–409.
- [9] C. Maya, S. Oruganti and R. Shivaji, Positive solutions for classes of  $p$ -Laplacian equations, *Dif. and Int. Eqns.*, 16(2003), 757–768.
- [10] N. N. Semenov, *Chemical Kinetics and Chain Reactions*, Oxford University Press, London, 1935.
- [11] R. Shivaji, A remark on the existence of three solutions via sub-super solutions. *Nonlinear Analysis and Application, Lecture Notes in Pure and Applied Mathematics*, 109(1987), 561–566.
- [12] Y. B. Zeldovich, G. I. Barenblatt, V. B. Librovich and G. M. Makhviladze, *The Mathematical Theory of Combustion and Explosions*, Consultants Bureau, New York, 1985.