# Existence Of Positive Solutions For A Class Of Kirchhoff Parabolic Systems With Multiple Parameters<sup>\*</sup>

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#### Abstract

In this paper, by using sub-super solutions method, we study the existence of weak positive solution for a class of Kirrchoff parabolic systems in bounded domains with multiple parameters. Our results are natural extensions from the previous ones in [2] and [8].

#### 1 Introduction

In this paper, we consider the following system of parabolic differential equations

$$\begin{pmatrix}
\frac{\partial u}{\partial t} - A\left(\int_{\Omega} |\nabla u|^2 dx\right) \Delta u = \lambda_1 \alpha(x) f(v) + \mu_1 \beta(x) h(u) \text{ in } Q_T, \\
\frac{\partial u}{\partial t} - B\left(\int_{\Omega} |\nabla v|^2 dx\right) \Delta v = \lambda_2 \gamma(x) g(u) + \mu_2 \eta(x) \tau(v) \text{ in } Q_T, \\
u = v = 0 \text{ on } \partial Q_T, \\
u(x, 0) = \varphi(x),
\end{cases}$$
(1)

where  $Q_T = \Omega \times [0,T]$ ,  $\Omega \subset \mathbb{R}^N$   $(N \geq 3)$  is a bounded smooth domain with  $C^2$  boundary  $\partial\Omega$ , and  $A, B : \mathbb{R}^+ \to \mathbb{R}^+$  are continuous functions,  $\alpha, \beta, \gamma, \eta \in C(\overline{\Omega})$ ,  $\lambda_1, \lambda_2, \mu_1$ , and  $\mu_2$  are nonnegative parameters.

Since the first equation in (1) contains an integral over  $\Omega$ , it is no longer a pointwise identity, therefore, it is often called nonlocal problem. This problem models several physical and biological systems, where u describes a process which depends on the

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average of itself, such as the population density, see [15]. Moreover, problem (1) is related to the stationary version of the Kirchhoff equation

$$\rho \frac{\partial^2 u}{\partial t^2} - \left(\frac{P_0}{h} + \frac{E}{2L} \int_0^L \left|\frac{\partial u}{\partial x}\right|^2 dx\right) \frac{\partial^2 u}{\partial x^2} = 0$$
(2)

presented by Kirchhoff in 1883 (see [11]). This equation is an extension of the classical d'Alembert's wave equation by considering the effect of the changes in the length of the string during the vibrations. The parameters in (2) have the following meanings: L is the length of the string, h is the area of the cross-section, E is the Young modulus of the material,  $\rho$  is the mass density, and  $P_0$  is the initial tension.

By using Euler time scheme on (1), we obtain the following problems

$$\begin{cases} u_{k} - \tau' A\left(\int_{\Omega} |\nabla u_{k}|^{2} dx\right) \Delta u = \tau' [\lambda_{1} \alpha (x) f(v) + \mu_{1} \beta (x) h(u_{k})] + u_{k-1} \text{ in } \Omega, \\ u_{k} - \tau' B\left(\int_{\Omega} |\nabla v|^{2} dx\right) \Delta v = \tau' [\lambda_{2} \gamma (x) g(u_{k}) + \mu_{2} \eta (x) \tau (v)] + u_{k-1} \text{ in } \Omega, \\ u_{k} = v = 0 \text{ on } \partial\Omega, \\ u_{0} = \rho, \end{cases}$$

$$(3)$$

where  $N\tau' = T$ ,  $0 < \tau' < 1$ , and for  $1 \le k \le N$ . In recent years, problems involving Kirchhoff type operators have been studied in many papers such as ([4], [12], [13], [16]–[18], [22]–[25]). In these articles, the authors have used different methods to get the existence of solutions for (1) in the single equation case. Z. Zhang in ([12] and [15]) studied the existence of nontrivial sign-changing solutions for system (1) where A(t) = B(t) = 1 via sub-supersolution method. Our paper is motivated by the recent results in [1], [2], [3], [8], [9] and [10]. Azzouz and Bensedik (Theorem 2 in [2]) investigated the existence of a positive solution for the nonlocal problem of the form

$$\begin{cases} -M\left(\int_{\Omega} |\nabla u|^2 \, dx\right) \Delta u = |u|^{p-2} \, u + \lambda f(x) \text{ in } \Omega,\\ u = 0 \text{ on } \partial\Omega, \end{cases}$$

$$\tag{4}$$

where  $\Omega$  is a bounded smooth domain in  $\mathbb{R}^N$ ,  $N \geq 3$  and p > 1, i.e. the nonlinear term at infinity and f is a sign-changing function. Using the sub and supersolution method combining a comparison principle introduced in [1], the authors established the existence of a positive solution for (4), where the parameter  $\lambda > 0$  is small enough. In the present paper, we consider system (1) in the case when the nonlinearities are "sublinear" at infinity, see the condition (H3). We are inspired by the ideas in the interesting paper [8], in which the authors considered system (1) in the case A(t) =B(t) = 1. More precisely, under suitable conditions on f, g, we shall show that system (1) has a positive solution for  $\lambda > \lambda^*$ . To our best knowledge, this is a new research topic for nonlocal problems (see [12, 13]). In the current paper, motivated by previous works in [2, 8] and by using sub-super solutions method, we study the existence of weak positive solution for a class of Kirchhoff parabolic systems in bounded domains with multiple parameters.

## 2 Existence Results

LEMMA 1 ([1]). Assume that  $M : \mathbb{R}^+ \to \mathbb{R}^+$  is a continuous and increasing function satisfying

$$\lim_{t \to 0^+} M\left(t\right) = m_0,\tag{5}$$

where  $m_0$  is a positive constant and assume that u, v are two non-negative functions such that

$$\begin{cases} -M\left(\int_{\Omega} |\nabla u|^2 dx\right) \Delta u \ge -M\left(\int_{\Omega} |\nabla v|^2 dx\right) \Delta v \text{ in } \Omega,\\ u = v = 0 \text{ on } \partial\Omega. \end{cases}$$
(6)

Then  $u \geq v$  in  $\Omega$ .

In this section, we shall state and prove the main result of this paper. Let us make the following assumptions:

(H1)  $A, B : \mathbb{R}^+ \to \mathbb{R}^+$  are two continuous and increasing functions and there exist  $a_i, b_i > 0, i = 1, 2$ , such that

$$a_1 \leq A(t) \leq a_2, \quad b_1 \leq B(t) \leq b_2 \text{ for all } t \in \mathbb{R}^+.$$

(H2)  $\alpha, \beta, \gamma, \eta \in C(\overline{\Omega})$  and

$$\alpha(x) \ge \alpha_0 > 0, \beta(x) \ge \beta_0 > 0, \gamma(x) \ge \gamma_0 > 0, \eta(x) \ge \eta_0 > 0,$$

for all  $x \in \Omega$ .

(H3)  $f, g, h, \text{and } \tau$  are continuous on  $[0, +\infty[, C^1 \text{ on } (0, +\infty)]$ , and increasing functions such that

$$\lim_{t \to +\infty} f(t) = +\infty, \quad \lim_{t \to +\infty} g(t) = +\infty, \quad \lim_{t \to +\infty} h(t) = +\infty = \lim_{t \to +\infty} \tau(t) = +\infty.$$

(H4) It holds that

$$\lim_{t \to +\infty} \frac{f\left(K\left(g\left(t\right)\right)\right)}{t} = 0, \text{ for all } K > 0$$

(H5)

$$\lim_{t \to +\infty} \frac{h(t)}{t} = \lim_{t \to +\infty} \frac{\tau(t)}{t} = 0.$$

THEOREM 1. Assume that the conditions (H1)–(H5) hold, and M is a nonincreasing function satisfying (5). Then for large  $\lambda_1 \alpha_0 + \mu_1 \beta_0$  and  $\lambda_2 \gamma_0 + \mu_2 \eta_0$ , problem (1) has a large positive weak solution.

We give the following two definitions before we prove our result.

DEFINITION 1.  $(u_k, v) \in (H_0^1(\Omega) \times H_0^1(\Omega))$  is said a weak solution of (3) if it satisfies

$$A\left(\int_{\Omega} |\nabla u_k|^2 dx\right) \int_{\Omega} \nabla u_k \nabla \phi dx = \int_{\Omega} \left[\lambda_1 \alpha\left(x\right) f\left(v\right) + \mu_1 \beta\left(x\right) h\left(u_k\right) - \frac{u_k - u_{k-1}}{\tau'}\right] \phi dx$$
$$B\left(\int_{\Omega} |\nabla v|^2 dx\right) \int_{\Omega} \nabla v \nabla \psi dx = \int_{\Omega} \left[\lambda_2 \gamma\left(x\right) g\left(u_k\right) \psi + \mu_2 \eta\left(x\right) \tau\left(v\right) - \frac{u_k - u_{k-1}}{\tau'}\right] \psi dx$$
for all  $(\phi, \psi) \in \left(H_0^1\left(\Omega\right) \times H_0^1\left(\Omega\right)\right)$ .

DEFINITION 2. A pair of nonnegative functions  $(\underline{u}_k, \underline{v})$  (respectively  $(\overline{u}_k, \overline{v})$ ) in  $(H_0^1(\Omega) \times H_0^1(\Omega))$  is called a weak subsolution (resp. supersolution pair) of (1) if they satisfy  $(u_k, \underline{v})$ ,  $(\overline{u}_k, \overline{v}) = (0, 0)$  on  $\partial\Omega$ ,

$$A\left(\int_{\Omega} \left|\nabla \underline{u_k}\right|^2 dx\right) \int_{\Omega} \nabla \underline{u_k} \nabla \phi dx \leq \int_{\Omega} \left[\lambda_1 \alpha\left(x\right) f\left(\underline{v}\right) + \mu_1 \beta\left(x\right) h\left(\underline{u_k}\right) - \frac{u_k - u_{k-1}}{\tau'}\right] \phi dx,$$

and

$$B\left(\int_{\Omega} \left|\nabla \underline{v}\right|^{2} dx\right) \int_{\Omega} \nabla \underline{v} \nabla \psi dx \leq \int_{\Omega} \left[\lambda_{2} \gamma\left(x\right) g\left(\underline{u_{k}}\right) + \mu_{2} \eta\left(x\right) \tau\left(\underline{v}\right) - \frac{u_{k} - u_{k-1}}{\tau'}\right] \psi dx$$

(respectively,

$$A\left(\int_{\Omega}\left|\nabla\overline{u_{k}}\right|^{2}dx\right)\int_{\Omega}\nabla\overline{u_{k}}\nabla\phi dx \geq \int_{\Omega}\left[\lambda_{1}\alpha\left(x\right)f\left(\overline{v}\right) + \mu_{1}\beta\left(x\right)h\left(\overline{u_{k}}\right) - \frac{u_{k} - u_{k-1}}{\tau'}\right]\phi dx$$

and

$$B\left(\int_{\Omega} |\nabla \overline{v}|^2 \, dx\right) \int_{\Omega} \nabla \overline{v} \nabla \psi \, dx \ge \int_{\Omega} \left[\lambda_2 \gamma\left(x\right) g\left(\overline{u_k}\right) + \mu_2 \eta\left(x\right) \tau\left(\overline{v}\right) - \frac{u_k - u_{k-1}}{\tau'}\right] \psi \, dx\right)$$

for all  $(\phi, \psi) \in (H_0^1(\Omega) \times H_0^1(\Omega))$ .

PROOF OF THEOREM 1. Let  $\sigma$  be the first eigenvalue of  $-\Delta$  with Dirichlet boundary conditions and  $\phi_1$  the corresponding positive eigenfunction with  $\|\phi_1\| = 1$ .

Let  $k_0, m_0, \delta > 0$  such that  $f(t), g(t), h(t), \tau(t) \ge -k_0$  for all  $t \in \mathbb{R}^+$  and  $|\nabla \phi_1|^2 - \sigma \phi_1^2 \ge m_0$  on  $\overline{\Omega}_{\delta} = \{x \in \Omega : d(x, \partial \Omega) \le \delta\}$ . For each  $\lambda_1 \alpha_0 + \mu_1 \beta_0$  and  $\lambda_2 \gamma_0 + \mu_2 \eta_0$  large, let us define

$$\underline{u_k} = \left(\frac{\left(\lambda_1 \alpha_0 + \mu_1 \beta_0\right) k_0}{2m_0 a_1}\right) \phi_1^2 \quad \text{and} \quad \underline{v} = \left(\frac{\left(\lambda_2 \gamma_0 + \mu_2 \eta_0\right) k_0}{2m_0 b_1}\right) \phi_1^2,$$

where  $a_1$  and  $b_1$  are given by the condition (H1). We shall verify that  $(\underline{u}_k, \underline{v})$  is a subsolution of problem (1) for  $\lambda_1 \alpha_0 + \mu_1 \beta_0$  and  $\lambda_2 \gamma_0 + \mu_2 \eta_0$  large enough. Indeed, let  $\phi \in H_0^1(\Omega)$  with  $\phi \ge 0$  in  $\Omega$ . By (H1)–(H3), a simple calculation shows that

$$\begin{split} &A\left(\int\limits_{\overline{\Omega}_{\delta}}\left|\nabla\underline{u}_{\underline{k}}\right|^{2}dx\right)\int\limits_{\overline{\Omega}_{\delta}}\nabla\underline{u}_{\underline{k}}.\nabla\phi dx\\ &= A\left(\int\limits_{\overline{\Omega}_{\delta}}\left|\nabla\underline{u}_{\underline{k}}\right|^{2}dx\right)\frac{(\lambda_{1}\alpha_{0}+\mu_{1}\beta_{0})\,k_{0}}{m_{0}a_{1}}\int\limits_{\overline{\Omega}_{\delta}}\phi_{1}\nabla\phi_{1}.\nabla\phi dx\\ &= \frac{(\lambda_{1}\alpha_{0}+\mu_{1}\beta_{0})\,k_{0}}{m_{0}a_{1}}A\left(\int\limits_{\overline{\Omega}_{\delta}}\left|\nabla\underline{u}_{\underline{k}}\right|^{2}dx\right)\left\{\int\limits_{\overline{\Omega}_{\delta}}\nabla\phi_{1}\nabla\left(\phi_{1}.\phi\right)dx-\int\limits_{\overline{\Omega}_{\delta}}\left|\nabla\phi_{1}\right|^{2}\phi dx\right\}\\ &= \frac{(\lambda_{1}\alpha_{0}+\mu_{1}\beta_{0})\,k_{0}}{m_{0}a_{1}}A\left(\int\limits_{\overline{\Omega}_{\delta}}\left|\nabla\underline{u}_{\underline{k}}\right|^{2}dx\right)\int\limits_{\overline{\Omega}_{\delta}}\left(\sigma\phi_{1}^{2}-\left|\nabla\phi_{1}\right|^{2}\right)\phi dx.\end{split}$$

On  $\overline{\Omega}_{\delta}$ , we have  $|\nabla \phi_1|^2 - \sigma \phi_1^2 \ge m_0$ . Then by using (H3),

$$f(\underline{v}), h(\underline{u_k}), g(\underline{u_k}), \tau(\underline{v}) \ge \frac{k_0}{m_0},$$

 $_{\rm thus}$ 

$$A\left(\int_{\overline{\Omega}_{\delta}} \left|\nabla \underline{u}_{\underline{k}}\right|^{2} dx\right) \int_{\overline{\Omega}_{\delta}} \nabla \underline{u}_{\underline{k}} \nabla \phi dx$$

$$\leq \frac{\left(\lambda_{1}\alpha_{0} + \mu_{1}\beta_{0}\right)k_{0}}{m_{0}} \int_{\overline{\Omega}_{\delta}} \left(\sigma\phi_{1}^{2} - \left|\nabla\phi_{1}\right|^{2}\right)\phi dx$$

$$\leq \int_{\Omega} \left[\lambda_{1}\alpha\left(x\right)f\left(\underline{v}\right) + \mu_{1}\beta\left(x\right)h\left(\underline{u}_{\underline{k}}\right) - \frac{u_{k} - u_{k-1}}{\tau'}\right]\phi dx.$$
(7)

Next, on  $\Omega \setminus \overline{\Omega}_{\delta}$ , we have  $\phi_1 \geq r$  for some r > 0. Therefore, under the conditions (H1)–(H3) and the definition of  $\underline{v}$ , it follows that

$$\int_{\Omega} \left[ \lambda_1 \alpha \left( x \right) f \left( \underline{v} \right) + \mu_1 \beta \left( x \right) h \left( \underline{u_k} \right) - \frac{u_k - u_{k-1}}{\tau'} \right] \phi \, dx$$

$$\geq (\lambda_{1}\alpha_{0} + \mu_{1}\beta_{0}) \frac{k_{0}a_{2}}{m_{0}a_{1}} \sigma \int_{\Omega \setminus \overline{\Omega}_{\delta}} \phi dx$$

$$\geq (\lambda_{1}\alpha_{0} + \mu_{1}\beta_{0}) \frac{k_{0}}{m_{0}a_{1}} A \left( \int_{\Omega \setminus \overline{\Omega}_{\delta}} |\nabla \underline{u}_{k}|^{2} dx \right) \sigma \int_{\Omega \setminus \overline{\Omega}_{\delta}} \phi dx$$

$$\geq (\lambda_{1}\alpha_{0} + \mu_{1}\beta_{0}) \frac{k_{0}}{m_{0}a_{1}} A \left( \int_{\Omega \setminus \overline{\Omega}_{\delta}} |\nabla \underline{u}_{k}|^{2} dx \right) \int_{\Omega \setminus \overline{\Omega}_{\delta}} (\sigma \phi_{1}^{2} - |\nabla \phi_{1}|^{2}) \phi dx$$

$$= A \left( \int_{\Omega \setminus \overline{\Omega}_{\delta}} |\nabla \underline{u}_{k}|^{2} dx \right) \int_{\Omega \setminus \overline{\Omega}_{\delta}} \nabla \underline{u}_{k} \nabla \phi dx, \qquad (8)$$

for  $\lambda_1 \alpha_0 + \mu_1 \beta_0 > 0$  large enough. Relations (7) and (8) imply that

$$A\left(\int_{\Omega} \left|\nabla \underline{u}_{k}\right|^{2} dx\right) \int_{\Omega} \nabla \underline{u}_{k} \nabla \phi dx$$
  
$$\leq \int_{\Omega} \left[\lambda_{1} \alpha\left(x\right) f\left(\underline{v}\right) + \mu_{1} \beta\left(x\right) h\left(\underline{u}_{k}\right) - \frac{u_{k} - u_{k-1}}{\tau'}\right] \phi dx, \qquad (9)$$

for  $\lambda_1 \alpha_0 + \mu_1 \beta_0 > 0$  large enough and any  $\phi \in H_0^1(\Omega)$  with  $\phi \ge 0$  in  $\Omega$ . Similarly,

$$B\left(\int_{\Omega} |\nabla \underline{v}|^{2} dx\right) \int_{\Omega} \nabla \underline{v} \nabla \psi dx$$
  

$$\leq \int_{\Omega} \left[\lambda_{2} \gamma\left(x\right) g\left(u_{k}\right) \psi + \mu_{2} \eta\left(x\right) \tau\left(v\right) - \frac{u_{k} - u_{k-1}}{\tau'}\right] \psi dx \qquad (10)$$

for  $\lambda_2 \gamma_0 + \mu_2 \eta_0 > 0$  large enough and any  $\psi \in H_0^1(\Omega)$  with  $\psi \ge 0$  in  $\Omega$ . From (9) and (10),  $(\underline{u}_k, \underline{v})$  is a subsolution of problem (3). Moreover, we have  $\underline{u}_k > 0, \underline{v} > 0$  in  $\Omega$ ,  $\underline{u} \to +\infty$  and  $\underline{v} \to +\infty$  also  $\lambda_1 \alpha_0 + \mu_1 \beta_0 \to +\infty$  and  $\lambda_2 \gamma_0 + \mu_2 \eta_0 \to +\infty$ . Next, we shall construct a supersolution of problem (3). Let  $\omega$  be the solution of

the following problem:

$$\begin{cases} -\triangle e = 1 \text{ in } \Omega, \\ e = 0 \text{ on } \partial\Omega. \end{cases}$$
(11)

Let

$$\overline{u_k} = Ce, \ \overline{v} = \left(\frac{\lambda_2 \left\|\gamma\right\|_{\infty} + \mu_2 \left\|\eta\right\|_{\infty}}{b_1}\right) \left[g\left(C \left\|e\right\|_{\infty}\right)\right]e,$$

where e is given by (11) and C > 0 is a large positive real number to be chosen later. We shall verify that  $(\overline{u_k}, \overline{v})$  is a supersolution of problem (3). Let  $\phi \in H_0^1(\Omega)$  with

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 $\phi \geq 0$  in  $\Omega.$  Then, we obtain from (11) and the condition (H1) that

$$\begin{split} A\left(\int_{\Omega} |\nabla \overline{u_k}|^2 \, dx\right) \int_{\Omega} \nabla \overline{u_k} \cdot \nabla \phi dx &= A\left(\int_{\Omega} |\nabla \overline{u_k}|^2 \, dx\right) C \int_{\Omega} \nabla \omega \cdot \nabla \phi dx \\ &= A\left(\int_{\Omega} |\nabla \overline{u_k}|^2 \, dx\right) C \int_{\Omega} \phi dx \\ &\geq a_1 C \int_{\Omega} \phi dx. \end{split}$$

By using (H4) and (H5), we can choose C large enough, thus

$$a_{1}C \geq \lambda_{1} \|\alpha\|_{\infty} f\left(\left[\frac{\lambda_{2} \|\gamma\|_{\infty} + \mu_{2} \|\eta\|_{\infty}}{b_{1}}\right] g\left(C \|e\|_{\infty}\right) \|e\|_{\infty}\right) + \mu_{1} \|\beta\|_{\infty} h\left(C \|e\|_{\infty}\right).$$

Therefore,

$$A\left(\int_{\Omega} |\nabla \overline{u_{k}}|^{2} dx\right) \int_{\Omega} \nabla \overline{u_{k}} \cdot \nabla \phi dx$$

$$\geq \left[\lambda_{1} \|\alpha\|_{\infty} f\left(\left[\frac{\lambda_{2} \|\gamma\|_{\infty} + \mu_{2} \|\eta\|_{\infty}}{b_{1}}\right] g\left(C \|e\|_{\infty}\right) \|e\|_{\infty}\right) + \mu_{1} \|\beta\|_{\infty} h\left(C \|e\|_{\infty}\right) \left] - \int_{\Omega} \frac{u_{k} - u_{k-1}}{\tau'} \phi dx$$

$$\geq \lambda_{1} \|\alpha\|_{\infty} \int_{\Omega} f\left(\left[\frac{\lambda_{2} \|\gamma\|_{\infty} + \mu_{2} \|\eta\|_{\infty}}{b_{1}}\right] g\left(C \|e\|_{\infty}\right) \|e\|_{\infty}\right) \phi dx$$

$$+ \mu_{1} \int_{\Omega} h\left(C \|e\|_{\infty}\right) \phi dx - \int_{\Omega} \frac{u_{k} - u_{k-1}}{\tau'} \phi dx$$

$$\geq \int_{\Omega} \left[\lambda_{1} \alpha\left(x\right) f\left(\underline{v}\right) + \mu_{1} \beta\left(x\right) h\left(\underline{u_{k}}\right) - \frac{u_{k} - u_{k-1}}{\tau'}\right] \phi dx.$$
(12)

Also, we have

$$B\left(\int_{\Omega} |\nabla \overline{v}|^{2} dx\right) \int_{\Omega} \nabla \overline{v} \nabla \psi dx$$
  

$$\geq (\lambda_{2} \|\gamma\|_{\infty} + \mu_{2} \|\eta\|_{\infty}) \int_{\Omega} g(C \|e\|_{\infty}) \psi dx$$
  

$$= \lambda_{2} \int_{\Omega} \gamma(x) g(\overline{u_{k}}) \psi dx + \mu_{2} \int_{\Omega} \eta(x) g(C \|e\|_{\infty}) \psi dx - \int_{\Omega} \frac{u_{k} - u_{k-1}}{\tau'} \psi dx.$$
(13)

Again by using (H4) and (H5) for C large enough, we have

$$g\left(C\left\|e\right\|_{\infty}\right) \ge \tau \left[\frac{\left(\lambda_{2}\left\|\gamma\right\|_{\infty} + \mu_{2}\left\|\eta\right\|_{\infty}\right)}{b_{1}}g\left(C\left\|e\right\|_{\infty}\right)\left\|e\right\|_{\infty}\right] \ge \tau\left(\overline{v}\right).$$
(14)

From (13) and (14), we have

$$B\left(\int_{\Omega} |\nabla \overline{v}|^{2} dx\right) \int_{\Omega} \nabla \overline{v} \nabla \psi dx$$
  

$$\geq \lambda_{2} \int_{\Omega} \gamma(x) g(\overline{u_{k}}) \psi dx + \mu_{2} \int_{\Omega} \eta(x) \tau(\overline{v}) \psi dx - \int_{\Omega} \frac{u_{k} - u_{k-1}}{\tau'} \psi dx. \quad (15)$$

From (12) and (15), we have  $(\overline{u}, \overline{v})$  is a subsolution of problem (1) with  $\underline{u} \leq \overline{u}$  and  $\underline{v} \leq \overline{v}$  for C large enough.

In order to obtain a weak solution of problem (3), we shall use the arguments by Azzouz and Bensedik [2] (observe that  $f, g, h, \text{ and } \tau$  does not depend on x). For this purpose, we define a sequence  $\{(u_n, v_n)\} \subset (H_0^1(\Omega) \times H_0^1(\Omega))$  as follows:  $u_0 = \overline{u}, v_0 = \overline{v}$  and  $(u_n, v_n)$  is the unique solution of the system

$$\begin{cases} -A\left(\int_{\Omega} |\nabla u_{n}|^{2} dx\right) \Delta u_{n} = \lambda_{1} \alpha\left(x\right) f\left(v_{n-1}\right) + \mu_{1} \beta\left(x\right) h\left(U_{n-1}\right) - \frac{u_{k} - u_{k-1}}{\tau'} \text{ in } \Omega, \\ -B\left(\int_{\Omega} |\nabla v_{n}|^{2} dx\right) \Delta v_{n} = \lambda_{2} \gamma\left(x\right) g\left(u_{n-1}\right) + \mu_{2} \eta\left(x\right) \tau\left(v_{n-1}\right) - \frac{u_{k} - u_{k-1}}{\tau'} \text{ in } \Omega, \\ u_{n} = v_{n} = 0 \text{ on } \partial\Omega. \end{cases}$$

$$(16)$$

We have  $(u_{n-1}, v_{n-1}) \in (H_0^1(\Omega) \times H_0^1(\Omega))$ , in the sense that, the right hand sides of (16) is independent on  $u_n$  and  $v_n$ . Setting

$$A(t) = tA(t^{2}), B(t) = tB(t^{2})$$

Since  $A(\mathbb{R}) = \mathbb{R}$ ,  $B(\mathbb{R}) = \mathbb{R}$ ,  $f(v_{n-1})$ ,  $h(u_{n-1})$ ,  $g(u_{n-1})$ , and  $\tau(v_{n-1}) \in L^2(\Omega)$ , we deduce from the results in [1], that system (16) has a unique solution  $(u_n, v_n) \in (H_0^1(\Omega) \times H_0^1(\Omega))$ . By using (16) and the fact that  $(u_0, v_0)$  is a supersolution of (1), we have

$$-A\left(\int_{\Omega} |\nabla u_{0}|^{2} dx\right) \Delta u_{0} \geq \lambda_{1} \alpha \left(x\right) f\left(v_{0}\right) + \mu_{1} \beta \left(x\right) h\left(u_{0}\right) - \frac{u_{k} - u_{k-1}}{\tau'}$$
$$= -A\left(\int_{\Omega} |\nabla u_{1}|^{2} dx\right) \Delta u_{1},$$
$$-B\left(\int_{\Omega} |\nabla v_{0}|^{2} dx\right) \Delta v_{0} \geq \lambda_{2} \gamma \left(x\right) g\left(u_{0}\right) + \mu_{2} \eta \left(x\right) \tau \left(v_{0}\right) - \frac{u_{k} - u_{k-1}}{\tau'}$$

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$$= -B\left(\int_{\Omega} |\nabla v_1| \, dx\right) \bigtriangleup v_1$$

and by using Lemma 1, we also have  $u_0 \ge u_1$  and  $v_0 \ge v_1$ . In addition, since  $u_0 \ge \underline{u}$ ,  $v_0 \ge \underline{v}$  and under the monotonicity condition of f, h, g, and  $\tau$ , we can deduce

$$-A\left(\int_{\Omega} |\nabla u_{1}|^{2} dx\right) \bigtriangleup u_{1} = \lambda_{1} \alpha(x) f(v_{0}) + \mu_{1} \beta(x) h(u_{0}) - \frac{u_{k} - u_{k-1}}{\tau'}$$
$$\geq \lambda_{1} \alpha(x) f(\underline{v}) + \mu_{1} \beta(x) h(\underline{u}) - \frac{u_{k} - u_{k-1}}{\tau'}$$
$$\geq -A\left(\int_{\Omega} |\nabla \underline{u}|^{2} dx\right) \bigtriangleup \underline{u}$$

 $\quad \text{and} \quad$ 

$$-B\left(\int_{\Omega} |\nabla v_{1}|^{2} dx\right) \Delta v_{1} = \lambda_{2} \gamma\left(x\right) g\left(u_{0}\right) + \mu_{2} \eta\left(x\right) \tau\left(v_{0}\right) - \frac{u_{k} - u_{k-1}}{\tau'}$$
$$\geq \lambda_{2} \gamma\left(x\right) g\left(\underline{u}\right) + \mu_{2} \eta\left(x\right) \tau\left(\underline{v}\right) - \frac{u_{k} - u_{k-1}}{\tau'}$$
$$\geq -B\left(\int_{\Omega} |\nabla \underline{v}|^{2} dx\right) \Delta \underline{v}.$$

According to Lemma 1, we have  $u_1 \geq \underline{u}, v_1 \geq \underline{v}$  for any  $u_2, v_2$ , thus we can write

$$-A\left(\int_{\Omega} |\nabla u_{1}|^{2} dx\right) \Delta u_{1} = \lambda_{1} \alpha(x) f(v_{0}) + \mu_{1} \beta(x) h(u_{0}) - \frac{u_{k} - u_{k-1}}{\tau'}$$
$$\geq \lambda_{1} \alpha(x) f(v_{1}) + \mu_{1} \beta(x) h(u_{0}) - \frac{u_{k} - u_{k-1}}{\tau'}$$
$$= -A\left(\int_{\Omega} |\nabla u_{2}|^{2} dx\right) \Delta u_{2},$$

$$-B\left(\int_{\Omega} |\nabla v_1| \, dx\right) \Delta v_1 = \lambda_2 \gamma\left(x\right) g\left(u_0\right) + \mu_2 \eta\left(x\right) \tau\left(v_0\right) - \frac{u_k - u_{k-1}}{\tau'}$$
$$\geq \lambda_1 \alpha\left(x\right) g\left(u_1\right) + \mu_2 \beta\left(x\right) \tau\left(v_1\right) - \frac{u_k - u_{k-1}}{\tau'}$$
$$= -B\left(\int_{\Omega} |\nabla v_2|^2 \, dx\right) \Delta v_2.$$

Then,  $u_1 \ge u_2, v_1 \ge v_2$ . Similarly,  $u_2 \ge \underline{u}$  and  $v_2 \ge \underline{v}$  because

$$-A\left(\int_{\Omega} |\nabla u_{2}|^{2} dx\right) \Delta u_{2} = \lambda_{1} \alpha(x) f(v_{1}) + \mu_{1} \beta(x) h(u_{1}) - \frac{u_{k} - u_{k-1}}{\tau'}$$

$$\geq \lambda_{1} \alpha(x) f(\underline{v}) + \mu_{1} \beta(x) h(\underline{u}) - \frac{u_{k} - u_{k-1}}{\tau'}$$

$$\geq -A\left(\int_{\Omega} |\nabla \underline{u}|^{2} dx\right) \Delta \underline{u},$$

$$-B\left(\int_{\Omega} |\nabla v_{2}|^{2} dx\right) \Delta v_{2} = \lambda_{2} \gamma\left(x\right) g\left(u_{1}\right) + \mu_{2} \eta\left(x\right) \tau\left(v_{1}\right) - \frac{u_{k} - u_{k-1}}{\tau'}$$
$$\geq \lambda_{2} \gamma\left(x\right) g\left(\underline{u}\right) + \mu_{2} \eta\left(x\right) \tau\left(\underline{v}\right) - \frac{u_{k} - u_{k-1}}{\tau'}$$
$$\geq -B\left(\int_{\Omega} |\nabla \underline{v}|^{2} dx\right) \Delta \underline{v}.$$

Repeating this argument, we get a bounded monotone sequence

$$\{(u_n, v_n)\} \subset \left(H_0^1(\Omega) \times H_0^1(\Omega)\right)$$

satisfying

$$\overline{u} = u_0 \ge u_1 \ge u_2 \ge \dots \ge u_n \ge \dots \ge \underline{u} > 0 \tag{17}$$

and

$$\overline{v} = v_0 \ge v_1 \ge v_2 \ge \dots \ge v_n \ge \dots \ge \underline{v} > 0.$$
(18)

Using the continuity of the functions  $f, h, g, \tau$  and the definition of the sequences  $\{u_n\}, \{v_n\}$ , there exist constants  $C_i > 0, i = 1, ..., 4$  independent of n such that

$$|f(v_{n-1})| \le C_1, \quad |h(u_{n-1})| \le C_2, \ |g(u_{n-1})| \le C_3$$
 (19)

and

$$|\tau(u_{n-1})| \leq C_4$$
 for all  $n$ .

Multiplying the first equation of (16) by  $u_n$ , integrating, using the Holder inequality and Sobolev embedding, we can show that

$$a_{1} \int_{\Omega} |\nabla u_{n}|^{2} dx \leq A \left( \int_{\Omega} |\nabla u_{n}|^{2} dx \right) \int_{\Omega} |\nabla u_{n}|^{2} dx$$
$$= \lambda_{1} \int_{\Omega} \alpha (x) f (v_{n-1}) u_{n} dx + \mu_{1} \int_{\Omega} \beta (x) h (u_{n-1}) u_{n} dx$$

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$$\begin{split} &-\int_{\Omega} \frac{u_{k} - u_{k-1}}{\tau'} u_{n} dx \\ &\leq \quad \lambda_{1} \left\| \alpha \right\|_{\infty} \int_{\Omega} |f(v_{n-1})| \ |u_{n}| \, dx + \mu_{1} \left\| \beta \right\|_{\infty} \int_{\Omega} |h(u_{n-1})| \ |u_{n}| \, dx \\ &- \int_{\Omega} \frac{u_{k} - u_{k-1}}{\tau'} \left| u_{n} \right| dx \\ &\leq \quad C_{1} \lambda_{1} \int_{\Omega} |u_{n}| \, dx + C_{2} \mu_{1} \int_{\Omega} |u_{n}| \, dx - \int_{\Omega} \frac{u_{k} - u_{k-1}}{\tau'} \left| u_{n} \right| dx \\ &\leq \quad C_{5} \left\| u_{n} \right\|_{H_{0}^{1}(\Omega)}, \end{split}$$

or

$$\|u_n\|_{H^1_0(\Omega)} \le C_5, \ \forall n,$$
 (20)

where  $C_5 > 0$  is a constant independent of n. Similarly, there exists  $C_6 > 0$  independent of n such that

$$\|v_n\|_{H^1_0(\Omega)} \le C_6, \quad \forall n. \tag{21}$$

From (20) and (21), we infer that  $\{(u_n, v_n)\}$  has a subsequence which weakly converges in  $H_0^1(\Omega)$  to a limit (u, v) with the properties  $u \ge \underline{u} > 0$  and  $v \ge \underline{v} > 0$ . Being monotone and also by using a standard regularity argument,  $\{(u_n, v_n)\}$  converges itself to (u, v).

Now, passing the limit in (16), we deduce that (u, v) is a positive solution of system (4). The proof of our theorem is completed.

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