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Global Existence And Uniqueness Of A Parabolic Haptotaxis Model^{*}

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Abstract

We study parabolic ODE systems modeling tumour invasion proposed by Anderson and Chaplain [3]. According to Yagi's arguments [12], we reduce them to corresponding evolution equations and show the existence of time global solutions.

1 Introduction

In this paper, we shall deal with the following parabolic system modeling haptotaxis

$$\partial_t u = D \Delta u - \rho \nabla . (u \nabla w),$$

$$\partial_t v = \delta \Delta v - \mu v + \alpha u, \qquad t > 0, \ x \in \Omega,$$

$$\partial_t w = -\gamma w v,$$

$$\begin{cases} \partial_n u = 0, \quad \partial_n v = 0, \quad \partial_n w_0 = 0 \quad \text{in } \partial\Omega, \\ u(0, .) = u_0, \ v(0, .) = v_0, \ w(0, .) = w_0 \quad \text{on } \Omega. \end{cases}$$
(1)

Here $\Omega \subset \mathbb{R}^d$ (d = 2, 3) is a bounded domain with C^3 boundary $\partial\Omega$ and the initial data u_0, v_0, w_0 are assumed to be nonnegative and ∂_n denotes the derivative with respect to the outer normal of $\partial\Omega$. This system is a mathematical model describing the motion of some species due to haptotaxis, the function u(t, x) corresponds to the cell density of the species at place $x \in \Omega$ and time $t \in [0, +\infty[$, and v(t, x) to the concentration of the chemical substance that is produced by the individuals while w = w(t, x) is the concentration of the extracellular matrix (ECM). The coefficients $D, \rho, \gamma, \delta, \alpha, \mu$ are given positive constants.

We first devote ourselves to the Cauchy problem for a semilinear evolution equation of the form (4) in a Banach space X. We present existence and uniqueness results in a way so that [12, Theorem 4.1] may be applied. Next, we use [12, Corollary 4.1] to show that the a priori estimate for local solutions of (4) with respect to the $A^{\frac{\beta}{2}}U(t)$

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norm ensures extension of local solutions without limit in order to construct the global solutions.

2 Local Existence of a Solution

Let Ω be a bounded domain in \mathbb{R}^d . For $1 \leq p \leq \infty$, $L^p(\Omega)$ is the usual Lebesgue space endowed with the norm $\|.\|_{L^p(\Omega)}$. Next for s > 0, $H^s(\Omega)$ is the usual fractional Sobolev space. We assume Ω has a C^3 class boundary $\partial\Omega$, and for $\frac{3}{2} < s \leq 3$

$$H_{N}^{s}\left(\Omega\right) = \left\{ u \in H^{s}\left(\Omega\right) : \ \partial_{n}u = 0 \text{ on } \partial\Omega \right\},\$$

and for $s < \frac{3}{2}$, we set $H_N^s(\Omega) = H^s(\Omega)$, with the norm $\|.\|_{H^s(\Omega)}$. We denote for $\frac{d}{2} < \beta < 2$ (d = 2, 3),

$$\mathcal{K} = \left\{ U_0 = (u_0, v_0, w_0)^t : \ 0 \le u_0 \in H_N^\beta(\Omega) , \ 0 \le v_0 \in H_N^{1+\beta}(\Omega) , \ 0 \le w_0 \in H_N^2(\Omega) \right\}.$$

The aim of this section is to prove the following Theorem:

THEOREM 2.1. Let β be a fixed exponent satisfying $d/2 < \beta < 2$ (d = 2, 3). For any $U_0 = (u_0, v_0, w_0) \in \mathcal{K}$, (1) possesses a unique local solution in the function space

$$u \in C([0, T_{U_0}]; H_N^2(\Omega)) \cap C([0, T_{U_0}]; H_N^\beta(\Omega)) \cap C^1([0, T_{U_0}]; L^2(\Omega)),$$

$$v \in C([0, T_{U_0}]; H_N^3(\Omega)) \cap C([0, T_{U_0}]; H_N^{1+\beta}(\Omega)) \cap C^1([0, T_{U_0}]; H^1(\Omega)),$$

$$w \in C([0, T_{U_0}]; H_N^2(\Omega)) \cap C^1([0, T_{U_0}]; H_N^2(\Omega)),$$
(2)

where $T_{U_0} > 0$ depends only on the norm $||u_0||_{H^{\beta}(\Omega)} + ||v_0||_{H^{\beta+1}(\Omega)} + ||w_0||_{H^2(\Omega)}$. In addition, for all $t \in [0, T_{U_0}]$, the solution satisfies the estimates

$$\|u(t)\|_{H^{\beta}(\Omega)} + \|v(t)\|_{H^{\beta+1}(\Omega)} + \|w(t)\|_{H^{2}(\Omega)} \le C_{U_{0}},$$
(3)

with some constant $C_{U_0} > 0$ depending on the norm $||u_0||_{H^{\beta}(\Omega)} + ||v_0||_{H^{\beta+1}(\Omega)} + ||w_0||_{H^2(\Omega)}$.

2.1 Proof of THEOREM 2.1

We formulate problem (1) as the Cauchy problem for an abstract semilinear equation

$$\begin{cases} \frac{dU}{dt} + AU = F(U), \\ U(0) = U_0, \end{cases}$$
(4)

in the Banach space

$$X = \left\{ U = (u, v, w)^{t} : \ u \in L^{2}(\Omega), \ v \in H^{1}(\Omega), \ w \in H^{2}_{N}(\Omega) \right\},\$$

endowed with the norm $||(u, v, w)^t|| = ||u||_{L^2(\Omega)} + ||v||_{H^1(\Omega)} + ||w||_{H^2(\Omega)}$ and A is a linear operator acting in X given by

$$A = diag \{A_1, A_2, A_3\} = diag \{-D\triangle + 1, -\delta\triangle + \mu, \gamma\}.$$

A is a sectorial linear operator of X, the spectrum of which is contained in a sectorial domain $\sigma(A) \subset \sum_{\omega} = \{\lambda \in \mathbb{C}, |\arg \lambda| < \omega_A\}$ with some angle $0 < \omega_A < \frac{\pi}{2}$. We refer to [12, Theorem 2.4] which ensures that the resolvent satisfies for $\lambda \notin \sigma(A)$ the estimate

$$\left\| \left(\lambda - A\right)^{-1} \right\| \leq \left\| \left(\lambda - A_1\right)^{-1} \right\|_{\mathcal{L}(L^2(\Omega))} + \left\| \left(\lambda - A_2\right)^{-1} \right\|_{\mathcal{L}(H^1(\Omega))} + \frac{1}{|\lambda - \gamma|} \\ \leq \frac{1 + \max\left\{ D, \frac{1}{D}, \frac{\delta}{\mu}, \frac{\mu}{\delta} \right\}}{|\lambda|}.$$

In $L_2(\Omega)$, under the Neumann boundary conditions on $\partial\Omega$, we have $\mathcal{D}(A_1) = H_N^2(\Omega)$ and according to [12, Theorem 16.7], we further have

$$\mathcal{D}(A_1^{\theta}) = \begin{cases} H^{2\theta}(\Omega), & 0 \le \theta < \frac{3}{4}, \\ H_N^{2\theta}(\Omega), & \frac{3}{4} < \theta \le 1, \end{cases}$$
(5)

with norm equivalence

$$c_{\Omega}^{-1} \|u\|_{H^{2\theta}(\Omega)} \le \left\|A_1^{\theta}u\right\|_{L^2(\Omega)} \le c_{\Omega} \|u\|_{H^{2\theta}(\Omega)}, \quad u \in \mathcal{D}(A_1^{\theta}).$$
(6)

In $H^1(\Omega)$, under the Neumann boundary conditions on $\partial\Omega$, it is known [12, Theorem 2.9] that $\mathcal{D}(A_2) = \{v \in H^2_N(\Omega) : \Delta v \in H^1(\Omega)\}$. Note that the fact that Ω has a C^3 class boundary $\partial\Omega$ ensures the shift property $\Delta v \in H^1(\Omega)$ with $\frac{\partial v}{\partial n} = 0$, implies that $\mathcal{D}(A_2) = H^3_N(\Omega)$; and according to [12, Theorem 16.1], we have $\mathcal{D}(A^2_{\theta}) =$ $[H^1(\Omega), H^3_N(\Omega)]_{\theta}, 0 \leq \theta \leq 1$. According to [12, Theorem 1.35],

$$\mathcal{D}(A_{2}^{\theta}) = \begin{cases} H^{2\theta+1}(\Omega), & 0 \le \theta < \frac{1}{4}, \\ H^{2\theta+1}_{N}(\Omega), & \frac{1}{4} < \theta \le 1, \end{cases}$$
(7)

with norm equivalence

$$c_{\Omega}^{-1} \|u\|_{H^{2\theta+1}(\Omega)} \le \left\|A_2^{\theta}u\right\|_{H^1(\Omega)} \le c_{\Omega} \|u\|_{H^{2\theta+1}(\Omega)}, \quad u \in \mathcal{D}(A_2^{\theta}),$$
(8)

where $c_{\Omega} > 0$ is determined by Ω . In $H_N^2(\Omega)$, the operator $A_3 = \gamma$ is a positive definite self-adjoint operator. By [12, Theorem 16.1] and [12, Theorem135], we have $\left[H_N^2(\Omega), H_N^2(\Omega)\right]_{\theta} = H_N^2(\Omega)$, therefore

$$\mathcal{D}(A_3^{\theta}) = H_N^2(\Omega) \ \ 0 \le \theta \le 1.$$
(9)

Consequently

$$\mathcal{D}(A) = \left\{ \left(u, v; w\right)^{t} : u \in H_{N}^{2}\left(\Omega\right), v \in H_{N}^{3}\left(\Omega\right), w \in H_{N}^{2}\left(\Omega\right) \right\}.$$
 (10)

Moreover it is clear that $A^{\theta} = diag \{A_1^{\theta}, A_2^{\theta}, A_3^{\theta}\}$. According to [12, Theorem 16.1], we have $\mathcal{D}(A^{\theta}) = [X, \mathcal{D}(A)]_{\theta}$. Then

$$\mathcal{D}(A^{\theta}) = \left\{ U = (u, v, w)^{t} \; ; \; u \in H^{2\theta}(\Omega) \; , \; v \in H^{2\theta+1}(\Omega) \; , \; w \in H^{2}_{N}(\Omega) \right\} \; , \; 0 < \theta < \frac{1}{4} \; ,$$
$$\mathcal{D}(A^{\theta}) = \left\{ U = (u, v, w)^{t} \; ; \; u \in H^{2\theta}(\Omega) \; , \; v \in H^{2\theta+1}_{N}(\Omega) \; , \; w \in H^{2}_{N}(\Omega) \right\} \; , \; \frac{1}{4} < \theta < \frac{3}{4} \; ,$$
$$\mathcal{D}(A^{\theta}) = \left\{ U = (u, v, w)^{t} \; ; \; u \in H^{2\theta}_{N}(\Omega) \; , \; v \in H^{2\theta+1}_{N}(\Omega) \; , \; w \in H^{2}_{N}(\Omega) \right\} \; , \; \frac{3}{4} < \theta \leq 1 \; .$$
(11)

The nonlinear operator F from $D(A^{\eta})$ ($\beta \leq \eta < 2$) into X is defined by

$$F(U) = (-\rho \nabla . (u \nabla w) + u, \ \alpha u, \ -\gamma (v-1)w)^t$$

Let $U, V \in D(A^{\eta})$ $(\beta \leq \eta < 2)$. Since $U = (u_1, v_1, w_1)^t$, $V = (u_2, v_2, w_2)^t$, we have

$$\|F(U) - F(V)\| \leq \rho \|\nabla (u_1 \nabla w_1 - u_2 \nabla w_2)\|_{L^2(\Omega)} + \alpha \|u_1 - u_2\|_{H^1(\Omega)} + \gamma \|(v_1 - 1)w_1 - (v_2 - 1)w_2\|_{H^2(\Omega)} + \|u_1 - u_2\|_{L^2(\Omega)}.$$
(12)

Since

 $\|\nabla . [u\nabla w]\|_{L^{2}(\Omega)} \leq \|\nabla u\|_{L^{4}(\Omega)} \|\nabla w\|_{L^{4}(\Omega)} + \|u\|_{L^{\infty}(\Omega)} \|\Delta w\|_{L^{2}(\Omega)},$

in the sequel, we need the following embeddings $H_N^\beta(\Omega) \to L^\infty(\Omega)$ and $H^1(\Omega) \to L^4(\Omega)$, $\frac{d}{2} < \beta < 2$ (d = 2, 3), to see that

$$\left\|\nabla \left[u \nabla w \right] \right\|_{L^{2}(\Omega)} \leq c_{\Omega} \left\| u \right\|_{H^{\beta}(\Omega)} \left\| w \right\|_{H^{2}(\Omega)}, \ u \in H^{\beta}(\Omega), w \in H^{2}_{N}(\Omega).$$
(13)

Moreover

$$\|\nabla. (u_1 \nabla w_1 - u_2 \nabla w_2)\|_{L^2(\Omega)}$$

 $\leq c_{\Omega} \|w_1\|_{H^2(\Omega)} \|u_1 - u_2\|_{H^{\beta}(\Omega)} + c_{\Omega} \|w_1 - w_2\|_{H^2(\Omega)} \|u_2\|_{H^{\beta}(\Omega)}.$ (14)

 $H^{2}\left(\Omega\right)$ is a Banach algebra, therefore

$$\| (v_1 - 1)w_1 - (v_2 - 1)w_2 \|_{H^2(\Omega)}$$

 $\leq c_{\Omega} \| w_1 - w_2 \|_{H^2(\Omega)} (\| v_1 \|_{H^2(\Omega)} + 1) + c_{\Omega} \| v_1 - v_2 \|_{H^2(\Omega)} \| w_2 \|_{H^2(\Omega)}.$ (15)

Let η be such that $\beta < \eta \leq 2$. Since $H_N^{\eta}(\Omega) \hookrightarrow H^1(\Omega) \hookrightarrow L^2(\Omega)$, we have

$$\|u_1 - u_2\|_{L^2(\Omega)} + \alpha \|u_1 - u_2\|_{H^1(\Omega)} \le c_\Omega (\alpha + 1) \|u_1 - u_2\|_{H^\eta(\Omega)}.$$
 (16)

We substitute (14), (15), (16) in (12) for $\beta \leq \eta < 2$. Then

$$\|F(U) - F(V)\| \le c_{\Omega} \left(\|u_{2}\|_{H^{\beta}(\Omega)} + \|v_{1}\|_{H^{\beta+1}(\Omega)} + 1 \right) \times \left[\|u_{1} - u_{2}\|_{H^{\eta}(\Omega)} + \|w_{1} - w_{2}\|_{H^{2}(\Omega)} + (\|w_{1}\|_{H^{2}(\Omega)} + \|w_{2}\|_{H^{2}(\Omega)}) \left(\|u_{1} - u_{2}\|_{H^{\beta}(\Omega)} + \|v_{1} - v_{2}\|_{H^{\beta+1}(\Omega)} \right) \right].$$

Therefore, in view of (11), (6), (8) and (9), we deduce that

$$\begin{aligned} \|F(U) - F(V)\| &\leq c_{\Omega} \left(\left\| A^{\frac{\beta}{2}}U \right\| + \left\| A^{\frac{\beta}{2}}V \right\| + 1 \right) \left[\left\| A^{\frac{\eta}{2}}(U - V) \right\| \\ &+ \left(\left\| A^{\frac{\eta}{2}}U \right\| + \left\| A^{\frac{\eta}{2}}V \right\| \right) \left\| A^{\frac{\beta}{2}}(U - V) \right\| \right], \ U, V \in D(A^{\eta}). \end{aligned}$$

Theorem 4.1 in [12] then provides the existence of local solutions. Indeed, for any $U_0 \in \mathcal{K}$, (4) possesses a unique local solution U in the function space:

$$U \in C((0, T_{U_0}]; \mathcal{D}(A)) \cap C([0, T_{U_0}]; \mathcal{D}(A^{\frac{p}{2}})) \cap C^1((0, T_{U_0}]; X).$$

Furthermore, the solution satisfies the estimates $||A^{\frac{\beta}{2}}U|| \leq C_{U_0}$. Here, $C_{U_0}, T_{U_0} > 0$ are determined by the norm $||U_0||_{\mathcal{D}(A^{\frac{\beta}{2}})}$ only. The proof of Theorem 2.1 is completed.

3 Nonnegativity of Local Solutions

We shall show that the local solution constructed above is nonnegative for $U_0 \in \mathcal{K}$. In the following we assume that $\Omega \subset \mathbb{R}^d (d = 2, 3)$ is a bounded domain with C^3 boundary. We denote by G the C^1 function defined by

$$G(s) = \begin{cases} \frac{1}{2}s^2, & s < 0, \\ 0, & s \ge 0. \end{cases}$$

PROPOSITION 3.1. Under the assumptions of THEOREM 2.1, we have

$$u(t,x) \ge 0, \ x \in \Omega, \ t \ge 0.$$

$$(17)$$

PROOF. We set $\psi(t) = \int_{\Omega} G(u(t,x)) dx$. We have $\psi'(t) = \int_{\Omega} G'(u) u_t dx$. Then

$$\psi'(t) = D \int_{\Omega} G'(u) \triangle u dx - \rho \int_{\Omega} G'(u) \nabla (u \nabla w) dx.$$

Observing that G'(u) = u if u < 0 and G'(u) = 0 if $u \ge 0$ and $G'(u) \in H^1(\Omega)$ for $u \in H^1(\Omega)$. Assuming $\frac{\partial w_0}{\partial n} = 0$ on $\partial\Omega$, we obtain $\frac{\partial w}{\partial n} = 0$ on $\partial\Omega$ and hence by Hölder's inequality, we have

$$\psi'(t) \le -D \|\nabla(G'(u))\|_{L^2(\Omega)}^2 + \frac{\rho}{2} \|G'(u)\|_{L^4(\Omega)}^2 \|\Delta w\|_{L^2(\Omega)}.$$
(18)

We use the interpolation inequality for d = 2, 3 to obtain

$$\|G'(u)\|_{L^4(\Omega)}^2 \le c_{\Omega} \|G'(u)\|_{H^1(\Omega)}^{\frac{d}{2}} \|G'(u)\|_{L^2(\Omega)}^{\frac{4-d}{2}}.$$

Then (3) shows that $\| \triangle w \|_{L^2(\Omega)} \leq C_{U_0}$, for $0 \leq t \leq T_{U_0}$. Therefore,

$$\rho \| \Delta w \|_{L^{2}(\Omega)} \| G'(u) \|_{L^{4}(\Omega)}^{2} \leq \frac{D}{2} \| \nabla (G'(u)) \|_{L^{2}(\Omega)}^{2} + C_{U_{0}} \| G'(u) \|_{L^{2}(\Omega)}^{2}.$$
(19)

Thus, in view of (19) and (18), $\psi'(t) \leq c_{T,U_0}\psi(t)$. By Gronwall's inequality $\psi(t) \leq \psi(0) \exp(tc_{T,U_0})$. Thus $\psi(0) = \int_{\Omega} G(u_0(t,x))dx = 0$ so that $\psi(t) = 0$. Hence $u \geq 0$.

PROPOSITION 3.2. Under the assumptions of Theorem 2.1, we have

$$v(t,x) \ge 0, \ x \in \Omega, \ t \ge 0.$$

$$(20)$$

PROOF. We set $\psi(t) = \int_{\Omega} G(v) dx$. Using the third equation of system (1), we have

$$\psi'(t) = -\delta \int_{\Omega} \left| \nabla G'(v) \right|^2 dx + \alpha \int_{\Omega} u G'(v) - \mu \int_{\Omega} v G'(v) dx,$$

since $vG'(v) \ge 0$, $G'(v) \le 0$, and $u \ge 0$ we have $\psi'(t) \le 0$, then $\psi(t) \le \psi(0)$. Since $\psi(0) = \int_{\Omega} G(v_0(t, x)) dx = 0$, we have $\psi(t) = 0$. Consequently $v \ge 0$.

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4 Global Solutions

In the following we assume that $\Omega \subset \mathbb{R}^d (d = 2, 3)$ is a bounded domain with C^3 boundary. As [12, Corollary 4.1] shows, the a priori estimates for local solutions of (4) with respect to the $A^{\frac{\beta}{2}}U(t)$ norm ensure extension of local solutions without limit. We may thus construct the global solutions.

For later use we state the following auxiliary results:

LEMMA 4.1. Under the assumptions of Theorem 2.1, for $0 \le t \le T_U$,

$$\|u(t)\|_{L^1(\Omega)} = \|u_0\|_{L^1(\Omega)}.$$
(21)

PROOF. Thanks to the homogeneous boundary conditions $\nabla u.\vec{n} = 0$ and $\nabla w_0.\vec{n} = 0$ on $\partial\Omega$, we may directly integrate (1) over Ω ; consequently $\int_{\Omega} \partial_t u dx = 0$, as $u \ge 0$, from (17) we have $\frac{d}{dt} ||u||_{L^1(\Omega)} = 0$, and mass conservation (21) is satisfied.

Next, we may easily prove the following lemma.

LEMMA 4.2. Suppose that $(u_0, v_0, w_0) \in \mathcal{K}$. Then for $0 \leq t \leq T_U$,

$$w(t,x) = w_0(x)e^{-\int_0^t \gamma v(\tau,x)d\tau}.$$
(22)

Moreover, we have

$$\|w(t)\|_{L^{\infty}(\Omega)} \le \|w_0\|_{L^{\infty}(\Omega)}.$$
(23)

PROPOSITION 4.3. Let Ω be a bounded smooth open domain of \mathbb{R}^d (d = 2, 3). Let $u \in H^1(\Omega)$. Then there exists a constant $c_{\Omega,\epsilon} > 0$ (depending on Ω, ϵ) such that

$$\|u\|_{L^{4}(\Omega)}^{2} \leq \epsilon \|\nabla u\|_{L^{2}(\Omega)}^{2} + c_{\Omega,\epsilon} \|u\|_{L^{1}(\Omega)}^{2}.$$
(24)

PROOF. With the help of the Cauchy inequality for the Gagliardo-Nirenberg inequality $||u||_{L^2(\Omega)}^2 \leq c_{\Omega} ||u||_{H^1(\Omega)}^{\frac{3d}{2+d}} ||u||_{L^1(\Omega)}^{\frac{4-d}{2+d}}$, (d = 2, 3), we get that

$$\|u\|_{L^{2}(\Omega)}^{2} \leq \frac{\epsilon}{4} \|\nabla u\|_{L^{2}(\Omega)}^{2} + \frac{1}{2} \|u\|_{L^{2}(\Omega)}^{2} + (\frac{c_{\Omega}^{2}}{2} + \frac{c_{\Omega}^{2}}{\epsilon}) \|u\|_{L^{1}(\Omega)}^{2}.$$
(25)

We simplify (25) so as to find

$$\|u\|_{L^{2}(\Omega)}^{2} \leq \frac{\epsilon}{2} \|\nabla u\|_{L^{2}(\Omega)}^{2} + \left(\frac{2c_{\Omega}^{2}}{\epsilon} + \frac{c_{\Omega}^{2}}{2}\right) \|u\|_{L^{1}(\Omega)}^{2}.$$
(26)

We take again the Gagliardo-Nirenberg's inequality $\|u\|_{L^4(\Omega)}^2 \leq c_{\Omega} \|u\|_{H^1(\Omega)}^{\frac{d}{2}} \cdot \|u\|_{L^2(\Omega)}^{\frac{4-d}{2}}$, with the Cauchy's inequality, then $\|u\|_{L^4(\Omega)}^2 \leq \frac{\epsilon}{2} \|\nabla u\|_{L^2(\Omega)} + (1 + \frac{c_{\Omega}^2}{2\epsilon}) \|u\|_{L^2(\Omega)}^2$. By combining with (26), (24) is proved. LEMMA 4.4. Suppose that $U_0 = (u_0, v_0, w_0) \in \mathcal{K}$. Then there exists a constant $c_{\Omega} > 0$ (depending on Ω) so that for $0 \le t \le T_U$,

$$\|v\|_{L^{2}(\Omega)} \leq c_{\Omega}(\|v_{0}\|_{L^{2}(\Omega)} + \|u_{0}\|_{L^{1}(\Omega)}).$$
(27)

PROOF. The second equation of (1) is written as the abstract equation

$$v(t) = e^{-tA_2}v_0 + \alpha \int_0^t e^{-(t-s)A_2}u \, ds, \ 0 \le t \le T_U,$$
(28)

in $L^2(\Omega)$. Therefore,

$$\begin{aligned} \|v\|_{L^{2}(\Omega)} &\leq \left\|e^{-tA_{2}}\right\|_{\mathcal{L}(L^{2}(\Omega))} \|v_{0}\|_{L^{2}(\Omega)} \\ &+ \alpha \int_{0}^{t} \left\|e^{-\frac{(t-s)}{2}A_{2}}\right\|_{\mathcal{L}(L^{2}(\Omega),L^{1}(\Omega))} \left\|e^{-\frac{(t-s)}{2}A_{2}}\right\|_{\mathcal{L}(L^{1}(\Omega))} \|u\|_{L^{1}(\Omega)} \, ds. \end{aligned}$$

From the estimate in [12, Eq. (2.128)] and [12, Theorem 2.28], and the formula $\mu^{-z}\Gamma(z) = \int_{0}^{+\infty} s^{z-1} e^{-\mu s} ds \ (Re(z) \in \mathbb{R}^{*}_{+}), \text{ we have, for } 0 \leq t \leq T_{U},$

$$\|v\|_{L^{2}(\Omega)} \leq c_{\Omega} \|v_{0}\|_{L^{2}(\Omega)} + \alpha c_{\Omega} \mu^{-\frac{4-d}{4}} \Gamma\left(\frac{4-d}{4}\right) \|u_{0}\|_{L^{1}(\Omega)}$$

The proof is completed.

We shall prove the following result.

PROPOSITION 4.5. Suppose that $U_0 = (u_0, v_0, w_0) \in \mathcal{K}$. Then

$$\left\|A^{\frac{\beta}{2}}U\right\| \le p(t + \|A^{\frac{\beta}{2}}U_0\|), \ 0 \le t \le T_U,\tag{29}$$

with some continuous increasing function p(.).

PROOF. We first derive the desired X bound. We employ a change of variable of the form $u \to \frac{u}{\varphi}$ where $\varphi(w) = e^{\frac{\rho}{D}w}$. This leads to the equation (1) in the form

$$\varphi(\frac{u}{\varphi})_t = D\nabla.\left(\varphi\nabla(\frac{u}{\varphi})\right) - u(\frac{\varphi_t}{\varphi}). \tag{30}$$

Moreover, φ satisfies $\varphi_t = \varphi'(w)w_t = -\frac{\gamma\rho}{D}\varphi wv$. By multiplying the equation (30) by $\frac{2u}{\varphi}$ and integrating over Ω , we obtain

$$\frac{d}{dt} \int_{\Omega} \varphi(\frac{u}{\varphi})^2 dx + D \int_{\Omega} \varphi \left| \nabla(\frac{u}{\varphi}) \right|^2 dx = \frac{\gamma \rho}{D} \int_{\Omega} (\frac{u}{\varphi})^2 \varphi w v dx.$$
(31)

Applying Hölder's inequality to (31),

$$\frac{d}{dt} \int_{\Omega} \varphi(\frac{u}{\varphi})^2 dx + D \int_{\Omega} \varphi \left| \nabla(\frac{u}{\varphi}) \right|^2 dx \le \frac{\gamma \rho}{D} \|\varphi w\|_{L^{\infty}(\Omega)} \|\frac{u}{\varphi}\|_{L^4(\Omega)}^2 \|v\|_{L^2(\Omega)},$$

in view of (23), we may then conclude that $\|\varphi w\|_{L^{\infty}(\Omega)} \leq \|w_0\|_{L^{\infty}(\Omega)} e^{\frac{\rho}{D} \|w_0\|_{L^{\infty}(\Omega)}}$. Hence, this, together with (27) and (24), yield that

$$\frac{d}{dt} \int_{\Omega} \varphi(\frac{u}{\varphi})^{2} dx + D \int_{\Omega} \varphi \left| \nabla(\frac{u}{\varphi}) \right|^{2} dx$$

$$\leq (\|v_{0}\|_{L^{2}(\Omega)} + \|u_{0}\|_{L^{2}(\Omega)}) \|w_{0}\|_{L^{\infty}(\Omega)} e^{\frac{\rho}{D} \|w_{0}\|_{L^{\infty}(\Omega)}} \|\frac{u}{\varphi}\|_{L^{4}(\Omega)}^{2}$$

$$\leq \frac{D}{2} \|\nabla(\frac{u}{\varphi})\|_{L^{2}(\Omega)}^{2} + \|\frac{u}{\varphi}\|_{L^{1}(\Omega)}^{2} c_{\|w_{0}\|_{H^{2}(\Omega)}, \|u_{0}\|_{L^{1}(\Omega)}, \|v_{0}\|_{L^{2}(\Omega)}}.$$
(32)

Since $(\frac{u}{\varphi})^2 \varphi = u^2 e^{-\frac{\rho}{D}w} \ge u^2 e^{-\frac{\rho}{D} ||w_0||_{L^{\infty}(\Omega)}}$, by solving the differential inequality (32), we obtain

$$\sup_{0 \le t \le T_U} \|u\|_{L^2(\Omega)}(t) \le tc(\|w_0\|_{L^{\infty}(\Omega)}, \|u_0\|_{L^2(\Omega)}, \|v_0\|_{L^2(\Omega)}).$$
(33)

It remains to prove the estimate in the space X for the two solution components v, w of (1). Multiplying the third equation of (1) by A_2v and integrating over Ω , and taking into account (8), we see that

$$\int_{0}^{t} \|v\|_{H^{2}(\Omega)}^{2} \leq c_{\Omega} \|v_{0}\|_{H^{1}(\Omega)}^{2} + c_{\Omega} \int_{0}^{t} \|u\|_{L^{2}(\Omega)}^{2} ds$$
$$\leq c_{\Omega} \|v_{0}\|_{H^{1}(\Omega)}^{2} + c_{\Omega} t \sup_{t \geq 0} \|u\|_{L^{2}(\Omega)}^{2}.$$
(34)

Next, we know that for $\Omega \subset \mathbb{R}^d$ (d = 2, 3), so

$$\|w_0 e^{-\int_0^t \gamma v}\|_{H^2(\Omega)}^2 \le c_\Omega \|w_0\|_{H^2(\Omega)}^2 \|e^{-\int_0^t \gamma v}\|_{H^2(\Omega)}^2.$$

Using the same arguments as in [12, inequality (13.34)], we see that

$$\|w\|_{H^{2}(\Omega)}^{2} \leq c_{\Omega} \|w_{0}\|_{H^{2}(\Omega)}^{2} \left(1 + \int_{0}^{t} \|v\|_{H^{2}(\Omega)}\right)^{2} \|e^{-\int_{0}^{t} \gamma v}\|_{L^{\infty}(\Omega)}^{2}.$$

Recalling (23), (34) and (33), we see that

$$\|w\|_{H^{2}(\Omega)} \leq (1+t^{2}) c(\|v_{0}\|_{H^{1}(\Omega)}, \|w_{0}\|_{H^{2}(\Omega)}, \|u_{0}\|_{L^{1}(\Omega)}, \Omega).$$
(35)

The next step is devoted to showing the estimate in the space $\mathcal{D}(A^{\frac{\beta}{2}})$. Using (13), [12, Eq. (2.128)] and [12, Theorem 2.28] with some exponent $\frac{d}{2} < \beta' < \beta$ and a constant

 $\begin{aligned} \left\| A_{1}^{\frac{\beta}{2}} u \right\|_{L^{2}(\Omega)} &\leq \left\| e^{-tA_{1}} \right\|_{\mathcal{L}(L^{2}(\Omega))} \left\| A_{1}^{\frac{\beta}{2}} u_{0} \right\|_{L^{2}(\Omega)} \\ &+ \int_{0}^{t} \left\| A_{1}^{\frac{\beta}{2}} e^{-\frac{1}{2}(t-s)A_{1}} \right\|_{\mathcal{L}(L^{2}(\Omega))} \left\| e^{-\frac{1}{2}(t-s)A_{1}} \right\|_{\mathcal{L}(L^{2}(\Omega))} \left\| \nabla . \left(u \nabla w \right) \right\|_{L^{2}(\Omega)} ds \\ &\leq c_{\Omega} \left\| u_{0} \right\|_{H^{\beta}(\Omega)} + \int_{0}^{t} c_{\Omega} \left(t-s \right)^{-\frac{\beta}{2}} e^{-\frac{1}{2}(t-s)} \left\| u \right\|_{H^{\beta'}(\Omega)} \left\| w \right\|_{H^{2}(\Omega)} ds. \end{aligned}$ (36)

By the same arguments as in [12, inequality (2.119)], (6), (33) and (35), we see that, for all $\varepsilon > 0$,

$$\begin{aligned} \|w\|_{H^{2}(\Omega)} \|u\|_{H^{\beta'}(\Omega)} &\leq c_{\Omega} \left(1+t^{2}\right) \left\|A_{1}^{\frac{\beta'}{2}}u\right\|_{L^{2}(\Omega)} \\ &\leq c_{\Omega} \left(1+t^{2}\right) \|u\|_{L^{2}(\Omega)}^{1-\frac{\beta'}{\beta}} \cdot \left\|A_{1}^{\frac{\beta}{2}}u\right\|_{L^{2}(\Omega)}^{\frac{\beta'}{\beta}} \\ &\leq c_{\Omega,\varepsilon} \left(1+t\right)^{\frac{3\beta}{\beta-\beta'}} + \varepsilon \left\|A_{1}^{\frac{\beta}{2}}u\right\|_{L^{2}(\Omega)}. \end{aligned}$$
(37)

Therefore, summing up (28), (37) and (36), we have for $0 \le t \le T_U$,

$$\begin{split} \sup_{0 \le t' \le t} \left\| A_1^{\frac{\beta}{2}} u \right\|_{L^2(\Omega)} & \le \quad c_\Omega \left\| u_0 \right\|_{H^{\beta}(\Omega)} + c_{\Omega,\varepsilon} \left(1 + t \right)^{\frac{3\beta}{\beta - \beta'}} \int_0^{+\infty} (t - s)^{-\frac{\beta}{2}} e^{-\frac{1}{2}(t - s)} ds \\ & + \frac{\varepsilon c_\Omega \int_0^{+\infty} (t - s)^{-\frac{\beta}{2}} e^{-\frac{1}{2}(t - s)} ds}{2} \sup_{0 \le t' \le t} \left\| A_1^{\frac{\beta}{2}} u \right\|_{L^2(\Omega)}. \end{split}$$

Let $\left(\frac{\beta}{2}\right)^{1-\frac{\beta}{2}}\Gamma\left(1-\frac{\beta}{2}\right) = \int_{0}^{+\infty} s^{-\frac{\beta}{2}} e^{-\frac{\beta s}{2}} ds$ and $\varepsilon^{-1} = c_{\Omega}\left(\frac{\beta}{2}\right)^{1-\frac{\beta}{2}}\Gamma\left(1-\frac{\beta}{2}\right)$. From (6) and (35), it follows that

$$\sup_{0 \le t \le T_U} \|u\|_{H^{\beta}(\Omega)} \le (1+t)^{\frac{3\beta}{\beta-\beta'}} c(\|v_0\|_{H^{1+\beta}(\Omega)}, \|w_0\|_{H^2(\Omega)}, \|u_0\|_{H^{\beta}(\Omega)}, \Omega).$$
(38)

 $c_{\Omega} > 0,$

In a similar manner, thanks to (28) and (38), we have for all $0 \le t \le T$,

$$\|v\|_{H^{1+\beta}(\Omega)} \leq \|v_0\|_{H^{1+\beta}(\Omega)} + \alpha c_{\Omega,\omega} \int_0^t \left\|A_2^{\frac{1}{2}} e^{-\frac{\varrho}{2}(t-s)A_2}\right\|_{\mathcal{L}(L^2(\Omega))} \left\|e^{-\frac{\varrho}{2}(t-s)A_1}\right\|_{\mathcal{L}(L^2(\Omega))} \|u\|_{H^{\beta}(\Omega)} \, ds$$

$$\leq \|v_0\|_{H^{1+\beta}(\Omega)} + \sup_{t\geq 0} \|u\|_{H^{\beta}(\Omega)} \, \alpha c_{\Omega,\omega} \int_0^t (t-s)^{-\frac{1}{2}} e^{-\frac{\varrho}{2}(t-s)} \, ds$$

$$\leq (1+t)^{\frac{3\beta}{\beta-\beta'}} \, c(\|v_0\|_{H^{1+\beta}(\Omega)}, \|v_0\|_{H^2(\Omega)}, \|u_0\|_{L^1(\Omega)}, \Omega). \tag{39}$$

Finally we use (35), (38) and (39), for $0 \le t \le T_U$, there exists a continuous increasing function $p(\cdot)$ such that

$$\|u\|_{H^{\beta}(\Omega)} + \|v\|_{H^{1+\beta}(\Omega)} + \|w\|_{H^{2}(\Omega)} \le p(t + \|u_{0}\|_{H^{\beta}(\Omega)} + \|v_{0}\|_{H^{1+\beta}(\Omega)} + \|w_{0}\|_{H^{2}(\Omega)}).$$

We will take the same steps and expressions of the proof of global existence in [12] to prove the following result.

THEOREM 4.6. For any $U_0 = (u_0, v_0, w_0) \in \mathcal{K}$, there exists a unique global solution of (1) in the function space:

$$\begin{aligned} u &\in C(]0, +\infty[; H_N^2(\Omega)) \cap C([0, +\infty[; H_N^\beta(\Omega)) \cap C^1]0, +\infty[; L^2(\Omega)), \\ v &\in C(]0, +\infty[; H_N^3(\Omega)) \cap C([0, +\infty[; H_N^{1+\beta}(\Omega)) \cap C^1(]0, +\infty[; H^1(\Omega)), \\ w &\in C([0, +\infty[; H_N^2(\Omega)) \cap C^1(]0, +\infty[; H_N^2(\Omega)). \end{aligned}$$

PROOF. Utilizing the a priori estimate (29), we shall construct a global solution to (1). For $U_0 \in \mathcal{K}$, we know that there exists a local solution at least on an interval $[0, T_{U_0}]$. Let $0 < t_1 < T_{U_0}$. Then, $U_1 = U(t_1) \in \mathcal{K}$. We next consider problem (1) with the initial value U_1 on an interval $[t_1, T]$, where the end time T > 0 is any finite time. The estimate (29) ensures for any local solution V, $\left\|A^{\frac{\beta}{2}}V(t)\right\| \leq p(\|A^{\frac{\beta}{2}}U_1\|+T)$, $t_1 \leq t \leq T_V$. Then, the local solution V can always be extended over an interval $[t_1, T_V + \tau]$ as local solution, $\tau > 0$ being dependent only on $p(\|A^{\frac{\beta}{2}}U_1\|_X + T)$ and hence being independent of the extreme time T_V (cf. [12. Corollary 4.1]). This means that our Cauchy problem possesses a global solution on the interval $[t_1, T]$.

This argument is meaningful for any finite time T > 0. So, we conclude the global existence of solution. For any initial value $U_0 \in \mathcal{K}$, there exists a unique global solution to (1) with $U(t) \in \mathcal{K}, 0 \leq t < \infty$, in the function space

$$U \in C(]0, +\infty[; \mathcal{D}(A)) \cap C([0, +\infty[; \mathcal{D}(A^{\frac{p}{2}}))) \cap C^{1}((]0, +\infty[; X).$$

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