

Three Applications In q^2 -Analogue Sobolev Spaces*

Ahmed Saoudi[†], Ahmed Fitouhi[‡]

Received 13 November 2015

Abstract

We study three q -difference-differential operators in q^2 -analogue Sobolev spaces.

1 Introduction

In this work, we are interested in the study of three q -difference-differential operators in q^2 -analogue Sobolev and q^2 -potential spaces introduced by the authors in their recent paper [8].

The outline of this paper is as follows. In section 2, we recall some basic facts from quantum calculus, some properties from the q -Rubin's operator and functional spaces in quantum calculus. We study in section 3 the hypoellipticity of q -Rubin operator. Section 4 is devoted to the existence and regularity of q^2 -analogue wave equation. In section 5, we study the solution of q^2 -analogue-Schrödinger equation.

2 Preliminary

Throughout this paper, we assume $0 < q < 1$ and we refer the reader to [3, 5] for the definitions and properties of hypergeometric functions. In this section we will fix some notations and recall some preliminary results. We put $\mathbb{R}_q = \{\pm q^n : n \in \mathbb{Z}\}$ and $\mathbb{R}_q = \mathbb{R}_q \cup \{0\}$. For $a \in \mathbb{C}$, the q -shifted factorials are defined by

$$(a; q)_0 = 1; \quad (a; q)_n = \prod_{k=0}^{n-1} (1 - aq^k), n = 1, 2, \dots; \quad (a; q)_\infty = \prod_{k=0}^{\infty} (1 - aq^k).$$

We denote also

$$[a]_q = \frac{1 - q^a}{1 - q}, \quad a \in \mathbb{C} \quad \text{and} \quad [n]_q! = \frac{(q; q)_n}{(1 - q)^n}, \quad n \in \mathbb{N}.$$

A q -analogue of the classical exponential function is given by (see [6, 7])

$$e(z; q^2) = \cos(-iz; q^2) + i \sin(-iz; q^2),$$

*Mathematics Subject Classifications: 35H10, 35L05.

[†]Department of Mathematics, Faculty of Sciences of Tunis, Tunis-El Manar University, Tunisia

[‡]Department of Mathematics, Faculty of Sciences of Tunis, Tunis-El Manar University, Tunisia

where

$$\cos(z; q^2) = \sum_{n=0}^{+\infty} q^{n(n+1)} \frac{(-1)^n z^{2n}}{[2n]q!} \quad \text{and} \quad \sin(z; q^2) = \sum_{n=0}^{+\infty} q^{n(n+1)} \frac{(-1)^n z^{2n+1}}{[2n+1]q!}.$$

The q -differential-difference operator is defined as (see [6, 7])

$$\partial_q f(z) = \begin{cases} \frac{f(q^{-1}z) + f(-q^{-1}z) - f(qz) + f(-qz) - 2f(-z)}{2(1-q)z} & \text{if } z \neq 0, \\ \lim_{z \rightarrow 0} \partial_q f(z) & \text{in } \mathbb{R}_q \quad \text{if } z = 0. \end{cases}$$

The q -Jackson integrals are defined by (see [4])

$$\int_0^a f(x) d_q x = (1-q)a \sum_{n=0}^{+\infty} q^n f(aq^n)$$

and

$$\int_{-\infty}^{+\infty} f(x) d_q x = (1-q) \sum_{n=-\infty}^{+\infty} q^n \{f(q^n) + f(-q^n)\},$$

provided the sums converge absolutely.

The q -Rubin-Fourier transform defined in [6], for all $x \in \tilde{\mathbb{R}}_q$ as

$$\mathcal{F}_q(f)(x) := K_q \int_{-\infty}^{+\infty} f(t) e(-itx; q^2) d_q t,$$

where

$$K_q = \frac{(1+q)^{\frac{1}{2}}}{2\Gamma_{q^2}(\frac{1}{2})}, \quad \text{and} \quad \Gamma_q(x) = \frac{(q; q)_\infty}{(q^x; q)_\infty} (1-q)^{1-x}.$$

In the following we denote by

- $\mathcal{S}_q(\mathbb{R}_q)$ stands for the q -analogue Schwartz space of smooth functions over \mathbb{R}_q whose q -derivatives of all order decay at infinity. $\mathcal{S}_q(\mathbb{R}_q)$ is endowed with the topology generated by the following family of semi-norms:

$$\|u\|_{n, \mathcal{S}_q}(f) := \sup_{x \in \mathbb{R}_q; k \leq n} (1+|x|)^n |\partial_q^k u(x)| \quad \text{for all } u \in \mathcal{S}_q \quad \text{and } n \in \mathbb{N}.$$

- $\mathcal{S}'_q(\mathbb{R}_q)$ the space of tempered distributions on \mathbb{R}_q , it is the topological dual of $\mathcal{S}_q(\mathbb{R}_q)$.
- $L_q^p(\mathbb{R}_q) = \left\{ f : \left(\int_{-\infty}^{+\infty} |f(x)|^p d_q x \right)^{\frac{1}{p}} < \infty \right\}$.
- $L_q^\infty(\mathbb{R}_q) = \left\{ f : \sup_{x \in I\mathbb{R}_q} |f(x)| < \infty \right\}$.
- $\mathcal{E}(\mathbb{R}; \mathcal{S}_q(\mathbb{R}_q))$ the space of C^∞ functions from \mathbb{R} into $\mathcal{S}_q(\mathbb{R}_q)$.

- $\mathcal{E}(\mathbb{R}; \mathcal{S}'_q(\mathbb{R}_q))$ the space of C^∞ functions from \mathbb{R} into $\mathcal{S}'_q(\mathbb{R}_q)$.
- $\mathcal{S}(\mathbb{R}; \mathcal{S}_q(\mathbb{R}_q))$ the space of Shawartz functions from \mathbb{R} into $\mathcal{S}_q(\mathbb{R}_q)$.
- $C(\mathbb{R}; \mathcal{W}_q^{s,p}(\mathbb{R}_q))$ the space of continuous functions on \mathbb{R} into $\mathcal{W}_q^{s,p}(\mathbb{R}_q)$.

It was shown in ([2, 6]) that \mathcal{F}_q verifies the following properties for $f \in L^1_q(\mathbb{R}_q)$.

1. If $uf(u) \in L^1_q(\mathbb{R}_q)$ and $\partial_q f \in L^1_q(\mathbb{R}_q)$ then

$$\partial_q(\mathcal{F}_q)(f)(x) = \mathcal{F}_q(-iuf(u))(x) \quad \text{and} \quad \mathcal{F}_q(\partial_q(f))(x) = ix\mathcal{F}_q(f)(x).$$

2. The reciprocity formula

$$\text{for } t \in \mathbb{R}_q, \quad f(t) = K \int_{-\infty}^{+\infty} \mathcal{F}_q(f)(x)e(itx; q^2)d_qx. \quad (1)$$

3. The q -Rubin-Fourier transform \mathcal{F}_q is an isomorphism from $\mathcal{S}_q(\mathbb{R}_q)$ onto itself and we have

$$\mathcal{F}_q^{-1}(f)(x) = \mathcal{F}_q(f)(-x) = \overline{\mathcal{F}_q(\bar{f})}(x) \quad \text{for } f \in \mathcal{S}_q(\mathbb{R}_q). \quad (2)$$

4. \mathcal{F}_q is an isomorphism from $L^2_q(\mathbb{R}_q)$ onto itself, and we have

$$\|\mathcal{F}_q(f)\|_{L^2_q(\mathbb{R}_q)} = \|f\|_{L^2_q(\mathbb{R}_q)} \quad \text{for } f \in L^2_q(\mathbb{R}_q). \quad (3)$$

DEFINITION 1. The q -Rubin-Fourier transform of a q -distribution $\nu \in \mathcal{S}'(\mathbb{R}_q)$ is defined by

$$\langle \mathcal{F}_q(\nu), \phi \rangle = \langle \nu, \mathcal{F}_q(\phi) \rangle \quad \text{for } \phi \in \mathcal{S}_q(\mathbb{R}_q). \quad (4)$$

From the above properties, we have the following result.

PROPOSITION 1. The q -Rubin-Fourier transform is a topological isomorphism from $\mathcal{S}'(\mathbb{R}_q)$ onto itself.

Let u be in $\mathcal{S}'(\mathbb{R}_q)$. We define the distribution $\partial_q u$, by

$$\langle \partial_q u, \psi \rangle = -\langle u, \partial_q \psi \rangle \quad \text{for } \psi \in \mathcal{S}_q(\mathbb{R}_q). \quad (5)$$

Hence, if we denote the q^2 -analogue Laplace operator by $\Delta_q := \partial_q^2$ we deduce

$$\langle \Delta_q u, \psi \rangle = \langle u, \Delta_q \psi \rangle \quad \text{for } \psi \in \mathcal{S}_q(\mathbb{R}_q). \quad (6)$$

These distributions satisfy the following properties

$$\mathcal{F}_q(\partial_q^n u) = (iy)^n \mathcal{F}_q(u), \quad n \in \mathbb{N}. \quad (7)$$

The q -translation operator $\tau_{q;x}$, $x \in \mathbb{R}_q$ is defined on $L_q^1(\mathbb{R}_q)$ by (see [6])

$$\tau_{q,y}(f)(x) = K \int_{-\infty}^{+\infty} \mathcal{F}_q(f)(t) e(itx; q^2) e(ity; q^2) d_q t, \quad y \in \mathbb{R}_q,$$

$$\tau_{q,0}(f)(x) = f(x).$$

DEFINITION 2. For $f \in L_q^2(\mathbb{R}_q)$ and $g \in L_q^1(\mathbb{R}_q)$, the q -convolution product is given by

$$f * g(y) = K \int_{-\infty}^{+\infty} \tau_{q,y} f(x) g(x) d_q x.$$

PROPOSITION 2. For $f \in L_q^n(\mathbb{R}_q)$, $g \in L_q^p(\mathbb{R}_q)$ and $1 \leq n, p, r \leq \infty$ such that $\frac{1}{n} + \frac{1}{p} - \frac{1}{r} = 1$ we have

$$\|f * g\|_{L_q^r(\mathbb{R}_q)} \leq \|f\|_{L_q^n(\mathbb{R}_q)} \|g\|_{L_q^p(\mathbb{R}_q)}. \quad (8)$$

DEFINITION 3. The q^2 -analogue Sobolev spaces introduced in [8] for $s \in \mathbb{R}$ and $1 \leq p < \infty$ are

$$\mathcal{W}_q^{s,p}(\mathbb{R}_q) := \{u \in \mathcal{S}'_q(\mathbb{R}_q) : (1 + |\xi|^2)^{\frac{s}{2}} \mathcal{F}_q(u) \in L_q^p(\mathbb{R}_q)\},$$

provided with the norm

$$\|u\|_{\mathcal{W}_q^{s,p}(\mathbb{R}_q)} := \left(\int_{-\infty}^{+\infty} (1 + \xi^2)^{\frac{sp}{2}} |\mathcal{F}_q(u)(\xi)|^p d_q \xi \right)^{\frac{1}{p}}.$$

DEFINITION 4. For $u \in \mathcal{S}'_q(\mathbb{R}_q)$ and $s \in \mathbb{R}$, the q^2 -potential operator \mathcal{P}_q^s of order s is defined in [1] as

$$\begin{aligned} \mathcal{P}_q^s : \mathcal{S}'_q(\mathbb{R}_q) &\longrightarrow \mathcal{S}'_q(\mathbb{R}_q) \\ u &\longmapsto (\mathcal{F}_q)^{-1}((1 + \xi^2)^{-\frac{s}{2}} \mathcal{F}_q(u)). \end{aligned}$$

DEFINITION 5. For all $(s, p) \in \mathbb{R} \times [1, +\infty[$, the q^2 -potential space is defined in [8] as

$$\mathcal{B}_q^{s,p}(\mathbb{R}_q) = \{u \in \mathcal{S}'_q(\mathbb{R}_q), \mathcal{P}_q^{-s}(u) \in L_q^p(\mathbb{R}_q)\},$$

provided with the norm

$$\|u\|_{\mathcal{B}_q^{s,p}(\mathbb{R}_q)} = \|\mathcal{P}_q^{-s}(u)\|_{L_q^p(\mathbb{R}_q)}.$$

3 Hypoellipticity of q -Rubin Operator

In this section, we will present the hypoellipticity of q -Rubin operator.

THEOREM 1. Let $P(\partial_q) = \sum_{j=0}^n \alpha_j (\partial_q)^j$, $\alpha_n \neq 0$, a q -differential-difference operator with constant coefficient α_j and symbol $P(\lambda) = \sum_{j=0}^n \alpha_j (i\lambda)^j \neq 0$, $\lambda \in \mathbb{R}$. If $u \in L_q^2(\mathbb{R}_q)$, $P(-\partial_q)u = f$ and $f \in L_q^2(\mathbb{R}_q)$, then $u \in \mathcal{B}_q^{n,2}(\mathbb{R}_q)$.

PROOF. One can easily show that there exists $R > 0$ and a positive constant C such that

$$|P(\xi)| \geq C|\xi|^n, \quad |\xi| \geq R. \quad (9)$$

Let $u \in \mathcal{S}_q(\mathbb{R}_q)$

$$\|u\|_{\mathcal{B}_q^{n,2}(\mathbb{R}_q)}^2 = \int_{\mathbb{R}_q} (\xi^2 + 1)^{n/2} |\mathcal{F}_q u(\xi)|^2 d_q \xi.$$

Taking $R \geq 1$, we have

$$\|u\|_{\mathcal{B}_q^{n,2}(\mathbb{R}_q)}^2 = \int_{-R}^R (\xi^2 + 1)^{n/2} |\mathcal{F}_q u(\xi)|^2 d_q \xi + \int_{|\xi| \geq R} (\xi^2 + 1)^{n/2} |\mathcal{F}_q u(\xi)|^2 d_q \xi.$$

Now, when $|\xi| \leq R$, we have $(\xi^2 + 1)^{n/2} \leq (R^2 + 1)^{n/2}$ and if $|\xi| \geq R$, $(\xi^2 + 1)^{n/2} \leq 2|\xi|^n$, we have

$$\|u\|_{\mathcal{B}_q^{n,2}(\mathbb{R}_q)}^2 \leq C(R^2 + 1)^{n/2} \int_{-R}^R |\mathcal{F}_q u(\xi)|^2 d_q \xi + C \int_{|\xi| \geq R} |\xi|^{2n} |\mathcal{F}_q u(\xi)|^2 d_q \xi.$$

According to the relations (3), (7) and (9) we obtain

$$\begin{aligned} \|u\|_{\mathcal{B}_q^{n,2}(\mathbb{R}_q)}^2 &\leq C \left(\int_{\mathbb{R}_q} |u(x)|^2 d_q x + \int_{\mathbb{R}_q} |\xi|^{2n} |\mathcal{F}_q u(\xi)|^2 d_q \xi \right) \\ &\leq C \left(\|u\|_{L_q^2(\mathbb{R}_q)}^2 + \int_{\mathbb{R}_q} |P(\xi) \mathcal{F}_q u(\xi)|^2 d_q \xi \right) \\ &\leq C \left(\|u\|_{L_q^2(\mathbb{R}_q)}^2 + \int_{\mathbb{R}_q} |\mathcal{F}_q (P(\partial_q)g)(\xi)|^2 d_q \xi \right) \\ &\leq C \left(\|u\|_{L_q^2(\mathbb{R}_q)}^2 + \|P(\partial_q)u\|_{L_q^2(\mathbb{R}_q)}^2 \right). \end{aligned}$$

The proof is completed by using the density of $\mathcal{S}_q(\mathbb{R}_q)$ in $\mathcal{B}_q^{n,2}(\mathbb{R}_q)$.

4 q^2 -Analogue Wave Equation

LEMMA 1. For all $p, n \in [1, \infty]$, $\frac{1}{r} = \frac{1}{p} + \frac{1}{n} - 1 \geq 0$, $s, s' \in \mathbb{R}$, $f \in \mathcal{B}_q^{s,p}(\mathbb{R}_q)$ and $g \in \mathcal{B}_q^{s',n}(\mathbb{R}_q)$, we have

$$f * g \in \mathcal{B}_q^{s+s',r}(\mathbb{R}_q) \quad \text{and} \quad \|f * g\|_{\mathcal{B}_q^{s+s',r}(\mathbb{R}_q)} \leq C_q \|f\|_{\mathcal{B}_q^{s,p}(\mathbb{R}_q)} \|g\|_{\mathcal{B}_q^{s',n}(\mathbb{R}_q)}.$$

PROOF. The results are given by the inequality (8) and the definition of the q^2 -potential spaces.

We consider the q^2 -analogue wave equation where the unknown u is a real-valued function such that

$$\begin{cases} \partial_t^2 u - \Delta_q u = 0', \\ u|_{t=0} = u_0 \in \mathcal{B}_q^{s,p}(\mathbb{R}_q)', \\ \partial_t u|_{t=0} = u_1 \in \mathcal{B}_q^{s',n}(\mathbb{R}_q)', \end{cases} \quad (t, x) \in \mathbb{R} \times \mathbb{R}_q.$$

COROLLARY 1. Let $\mathcal{C} := \{\xi \in \mathbb{R}_q, r \leq |\xi| \leq R\}$ for some positive reals r and R such that $r < R$. We assume that u_0 and u_1 are two functions satisfying

$$\text{supp } \mathcal{F}_q(u_j) \subset \mathcal{C}.$$

1. For $p = n = 2$, $u \in \mathcal{B}_q^{a+s,\infty}(\mathbb{R}_q) + \mathcal{B}_q^{b+s',\infty}(\mathbb{R}_q)$. For $a + s = b + s' = c$,

$$\|u\|_{\mathcal{B}_q^{c,\infty}(\mathbb{R}_q)} \leq C \left(\|u_0\|_{\mathcal{B}_q^{s,2}(\mathbb{R}_q)} + \|u_1\|_{\mathcal{B}_q^{s',2}(\mathbb{R}_q)} \right).$$

2. For $p \neq 2$ and $n \neq 2$, $u \in \mathcal{B}_q^{a+s, \frac{p}{2-p}}(\mathbb{R}_q) + \mathcal{B}_q^{b+s', \frac{n}{2-n}}(\mathbb{R}_q)$. For $a + s = b + s' = c$,

$$\|u\|_{\mathcal{B}_q^{c, \frac{p}{2-p}}(\mathbb{R}_q)} \leq C \left(\|u_0\|_{\mathcal{B}_q^{s,p}(\mathbb{R}_q)} + \|u_1\|_{\mathcal{B}_q^{s',p}(\mathbb{R}_q)} \right).$$

PROOF. According to the Duhamel expression for the solution and Lemma 1 we obtain the results.

5 q^2 -Analogue-Schrödinger Equation

Now we consider the following equation where the unknown u is a complex-valued function

$$\begin{cases} \partial_t u - i\Delta_q u = 0, \\ u|_{t=0} = g, \end{cases} \quad (t, x) \in \mathbb{R} \times \mathbb{R}_q.$$

THEOREM 2. Let $g \in \mathcal{S}'_q(\mathbb{R}_q)$. There exists a unique solution $u \in \mathcal{E}(\mathbb{R}; \mathcal{S}'_q(\mathbb{R}_q))$ such that

$$\begin{cases} \partial_t u - i\Delta_q u = 0, & \text{in } D'(\mathbb{R} \times \mathbb{R}_q), \\ u|_{t=0} = g. \end{cases}$$

PROOF. Let us prove the existence first. For $t \in \mathbb{R}$, we write

$$u_t = (\mathcal{F}_q)^{-1}(e^{-it|\xi|^2} \mathcal{F}_q(g)). \quad (10)$$

According to (4) we have

$$\langle u_t, \phi \rangle = \langle \mathcal{F}_q(g), e^{-it|\xi|^2} (\mathcal{F}_q)^{-1}(\phi) \rangle.$$

Therefore we deduce that $u_t \in \mathcal{E}(\mathbb{R}; \mathcal{S}'_q(\mathbb{R}_q))$, and $\mathcal{F}_q(u_t) \in \mathcal{E}(\mathbb{R}; \mathcal{S}'_q(\mathbb{R}_q))$. We recall that u is defined by

$$\langle u, \psi \rangle = \int_{\mathbb{R}} \langle u_t, \psi(t, \cdot) \rangle dt, \quad \psi \in \mathcal{S}(\mathbb{R}; \mathcal{S}_q(\mathbb{R}_q)).$$

Then, using (6), we have for any ψ in $\mathcal{S}(\mathbb{R}; \mathcal{S}_q(\mathbb{R}_q))$

$$\begin{aligned} \langle \partial_t u - i\Delta_q u, \psi \rangle &= -\langle u, \partial_t \psi(t, \cdot) + i\Delta_q \psi(t, \cdot) \rangle \\ &= -\int_{\mathbb{R}} \langle u_t, \partial_t \psi(t, \cdot) + i\Delta_q \psi(t, \cdot) \rangle dt \\ &= -\int_{\mathbb{R}} \langle \mathcal{F}_q(u_t), (\mathcal{F}_q)^{-1}(\partial_t \psi(t, \cdot) + i\Delta_q \psi(t, \cdot)) \rangle dt \\ &= -\int_{\mathbb{R}} \langle e^{-it|\cdot|^2} \mathcal{F}_q(g), (\partial_t - i|\cdot|^2)(\mathcal{F}_q)^{-1} \psi(t, \cdot) \rangle dt. \end{aligned}$$

Since

$$\partial_t \left(e^{-it|\xi|^2} (\mathcal{F}_q)^{-1} \psi(t, \xi) \right) = [\partial_t - i|\xi|^2 \mathcal{F}_q]^{-1} \psi(t, \xi) e^{-it|\xi|^2},$$

we see that

$$\begin{aligned} \langle \partial_t u + i\Delta_q u, \psi \rangle &= -\int_{\mathbb{R}} \langle \mathcal{F}_q(g), \partial_t \left(e^{-it|\cdot|^2} (\mathcal{F}_q)^{-1} \psi(t, \cdot) \right) \rangle dt \\ &= -\int_{\mathbb{R}} \partial_t \langle \mathcal{F}_q(g), e^{-it|\cdot|^2} (\mathcal{F}_q)^{-1} \psi(t, \cdot) \rangle dt = 0. \end{aligned}$$

Hence the existence of a solution u is shown. Let us now prove the uniqueness, which is equivalent to show that $u \equiv 0$ is the solution of the following problem

$$\begin{cases} \partial_t u - i\Delta_q u = 0 & \text{in } \mathcal{E}(\mathbb{R}; \mathcal{S}'_q(\mathbb{R}_q)), \\ u_{t=0} = 0. \end{cases}$$

In fact, for all ψ in $\mathcal{S}(\mathbb{R}; \mathcal{S}_q(\mathbb{R}_q))$ we have

$$\langle \partial_t u - i\Delta_q u, \psi \rangle = -\int_{\mathbb{R}} \langle u_t, (\partial_t + i\Delta_q) \psi(t, \cdot) \rangle dt = 0.$$

Although

$$\frac{d}{dt} \langle u_t, \psi(t, \cdot) \rangle = \langle u_t^{(1)}, \psi(t, \cdot) \rangle + \langle u_t, \partial_t \psi(t, \cdot) \rangle,$$

therefrom

$$-\int_{\mathbb{R}} \frac{d}{dt} \langle u_t, \psi(t, \cdot) \rangle dt + \int_{\mathbb{R}} \left[\langle u_t^{(1)}, \psi(t, \cdot) \rangle - i \langle u_t, \Delta_q \psi(t, \cdot) \rangle \right] dt = 0. \quad (11)$$

Since $\psi(-\infty, \cdot) = \psi(\infty, \cdot)$, we obtain

$$\int_{\mathbb{R}} \left[\langle u_t^{(1)}, \psi(t, \cdot) \rangle - i \langle u_t, \Delta_q \psi(t, \cdot) \rangle \right] dt = 0. \quad (12)$$

Besides, using the fact that $\mathcal{F}_q(u_t^{(1)}) = (\mathcal{F}_q(u_t))^{(1)}$ and the relations (7) and (12) we deduce that

$$\int_{\mathbb{R}} \left[\langle (\mathcal{F}_q(u_t))^{(1)}, (\mathcal{F}_q)^{-1} \psi(t, \cdot) \rangle + i \langle \mathcal{F}_q(u_t), |\cdot|^2 (\mathcal{F}_q)^{-1} \psi(t, \cdot) \rangle \right] dt = 0 \quad (13)$$

for $\psi \in \mathcal{S}(\mathbb{R}; \mathcal{S}_q(\mathbb{R}_q))$. If we choose ψ such that $(\mathcal{F}_q)^{-1} \psi(t, \xi) = e^{it|\xi|^2} \varphi(\xi) \chi(t)$ where φ in $\mathcal{S}_q(\mathbb{R}_q)$, χ in $\mathcal{S}(\mathbb{R})$, we have

$$\int_{\mathbb{R}} \left[\langle (\mathcal{F}_q(u_t))^{(1)}, e^{it|\cdot|^2} \varphi \rangle + i \langle \mathcal{F}_q(u_t), |\cdot|^2 e^{it|\cdot|^2} \varphi \rangle \right] \chi(t) dt = 0 \quad (14)$$

for $\chi \in \mathcal{S}(\mathbb{R})$. Hence we deduce that

$$\frac{d}{dt} \langle \mathcal{F}_q(u_t), e^{it|\cdot|^2} \varphi \rangle = \langle (\mathcal{F}_q(u_t))^{(1)}, e^{it|\cdot|^2} \varphi \rangle + i \langle \mathcal{F}_q(u_t), |\cdot|^2 e^{it|\cdot|^2} \varphi \rangle = 0 \quad (15)$$

for $\varphi \in \mathcal{S}_q(\mathbb{R}_q)$. Thus for all φ in $\mathcal{S}_q(\mathbb{R}_q)$, the function $t \mapsto \langle \mathcal{F}_q(u_t), e^{it|\cdot|^2} \varphi \rangle$ is constant. Finally, as $u_0 = 0$ then

$$\langle \mathcal{F}_q(u_t), e^{it|\cdot|^2} \varphi \rangle = \langle \mathcal{F}_q(u_0), \varphi \rangle = 0 \quad \text{for } t \in \mathbb{R} \text{ and } \varphi \in \mathcal{S}_q(\mathbb{R}_q).$$

Then we deduce that $u = 0$.

THEOREM 3. Let g be in $\mathcal{W}_q^{s,p}(\mathbb{R}_q)$, $s \in \mathbb{R}$ and $1 \leq p < \infty$, the solution given by the Theorem 2 belongs to $C(\mathbb{R}; \mathcal{W}_q^{s,p}(\mathbb{R}_q))$. For $m \in \mathbb{N}$, $(u_t^{(m)}) \in C(\mathbb{R}; \mathcal{W}_q^{s-2m,p}(\mathbb{R}_q))$ and we have for $t \in \mathbb{R}$

$$\begin{cases} \|u_t\|_{\mathcal{W}_q^{s,p}(\mathbb{R}_q)} = \|g\|_{\mathcal{W}_q^{s,p}(\mathbb{R}_q)}, \\ \|u_t^{(m)}\|_{\mathcal{W}_q^{s-2m,p}(\mathbb{R}_q)} \leq C_m \|g\|_{\mathcal{W}_q^{s,p}(\mathbb{R}_q)} \quad \text{for } m \in \mathbb{N}^*. \end{cases} \quad (16)$$

PROOF. By the formula (10), we have for all t in \mathbb{R}

$$\mathcal{F}_q(u_t) = e^{-it|\xi|^2} \mathcal{F}_q(g),$$

so, it is easy to deduce (16). Now, we will prove that for $m = 1, 2, \dots, u_t^{(m)}$ belongs to $C(\mathbb{R}; \mathcal{W}_q^{s-2m,p}(\mathbb{R}_q))$. In fact, let $(t_n)_n$ be a sequence that converge to t_0 in \mathbb{R} , we have

$$\|u_{t_n} - u_{t_0}\|_{\mathcal{W}_q^{s,p}(\mathbb{R}_q)}^2 = \int_{\mathbb{R}_q} (1 + |\xi|^2)^{\frac{sp}{2}} |e^{-it_n|\xi|^2} - e^{-it_0|\xi|^2}|^p |\mathcal{F}_q(g)(\xi)|^p d_q \xi.$$

According to the dominated convergence theorem, we obtain

$$\lim_{n \rightarrow \infty} \|u_{t_n} - u_{t_0}\|_{\mathcal{W}_q^{s,p}(\mathbb{R}_q)}^2 = 0.$$

Elsewhere, from (10) we have

$$\mathcal{F}_q(u_t^{(m)}) = (-i|\xi|^2)^m e^{-it|\xi|^2} \mathcal{F}_q(g).$$

Hence, we obtain

$$\|u_{t_n}^{(m)} - u_{t_0}^{(m)}\|_{\mathcal{W}_q^{s,p}(\mathbb{R}_q)}^p = \int_{\mathbb{R}_q} (1 + |\xi|^2)^{\frac{sp}{2}} |e^{-it_n|\xi|^2} - e^{-it_0|\xi|^2}|^p |\xi|^{2mp} |\mathcal{F}_q(g)(\xi)|^p d_q \xi.$$

Finally, the dominated convergence theorem leads to the result.

References

- [1] N. Bettaibi, M. M. Chaffar and A. Fitouhi, Sobolev type spaces associated with the q -Rubin's operator, *Le Matematiche*, 69(2014), 36–56.
- [2] A. Fitouhi and R. H. Bettaieb, Wavelet transforms in the q^2 -analogue Fourier analysis, *Math. Sci. Res. J.*, 12(2008), 202–214.
- [3] G. Gasper and M. Rahman, *Basic Hypergeometric Series*, vol. 96, Cambridge university press, 2004.
- [4] F. Jackson, On a q -definite integrals, *Quart. J. Pure Appl. Math.*, 41(1910), 193–203.
- [5] V. Kac and P. Cheung, *Quantum Calculus*, Springer-Verlag, New York, 2002.
- [6] R. L. Rubin, Duhamel solutions of non-homogeneous q^2 -analogue wave equations, *Proc. Amer. Math. Soc.*, 135(2007), 777–785.
- [7] R. L. Rubin, A q^2 -analogue operator for q^2 -analogue Fourier analysis, *J. Math. Anal. Appl.*, 212(1997), 571–582.
- [8] A. Saoudi and A. Fitouhi, On q^2 -analogue Sobolev type spaces, *Le Matematiche*, 70(2015), 63–77.