

New Exact Traveling Wave Solutions for Compound KdV-Burgers Equations in Mathematical Physics ^{*†}

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Abstract

With the aid of Mathematica, new explicit and exact travelling wave and solitary solutions for compound KdV-Burgers equations are obtained by using an improved sine-cosine method and the Wu elimination method.

1 Introduction

In the present paper we consider the compound KdV-Burgers equation

$$u_t + puu_x + qu^2u_x + ru_{xx} - su_{xxx} = 0, \quad (1)$$

where p, q, r, s are constants. This equation can be thought of as a generalization of the KdV, mKdV and Burgers equations, involving nonlinear dispersion and dissipation effects. As particular cases, (i) when $r = 0$ and $pqs \neq 0$, (1) becomes the compound KdV equation

$$u_t + puu_x + qu^2u_x - su_{xxx} = 0, \quad (2)$$

(ii) when $p = 0$ and $qrs \neq 0$, (1) becomes the KdV-Burgers equation

$$u_t + qu^2u_x + ru_{xx} - su_{xxx} = 0, \quad (3)$$

and (iii), when $r = 0$ in (3), then we get the mKdV equation

$$u_t + qu^2u_x - su_{xxx} = 0. \quad (4)$$

In a recent paper, Wang [1] has found some exact solutions of (1) by using the homogenous balance method. In this paper we obtain new travelling wave solutions of (1) by using an improved sine-cosine method [2,3] and Wu's elimination method [4]. The main idea of the algorithm is as follows. Given a partial differential equation of the form

$$P(u, u_x, u_t, u_{xx}, u_{xt}, u_{tt}, \dots) = 0, \quad (5)$$

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where P is a polynomial. By assuming travelling wave solutions of the form

$$u(x, t) = \varphi(\xi), \xi = \lambda(x - kt + c_0), \quad (6)$$

where k, λ are constants to be determined, and c_0 is an arbitrary constant, we are led to the ordinary differential equation

$$P(\varphi, \varphi', \varphi'', \dots) = 0, \quad (7)$$

where φ' denotes $d\varphi/d\xi$. According to the sine-cosine method (see [1-6] for details), we suppose that (7) has the following formal travelling wave solution

$$\varphi(\xi) = \sum_{i=1}^n \sin^{i-1} \omega (B_i \sin \omega + A_i \cos \omega) + A_0, \quad (8)$$

and

$$\frac{d\omega}{d\xi} = \sin \omega, \text{ or } \frac{d\omega}{d\xi} = \cos \omega, \quad (9)$$

where A_0, \dots, A_n and B_1, \dots, B_n are constants to be determined. Then we proceed as follows:

Step 1. Equating the highest order nonlinear term and highest order linear partial derivative in (5), yield the value of n .

Step 2. Substituting (8), (9) into (7), we obtain a polynomial equation involving $\cos \omega \sin^i \omega, \sin^i \omega$ for $i = 0, 1, 2, \dots, n$. This step can be carried out the help of Mathematica.

Step 3. Setting the constant term and coefficients of $\sin \omega, \cos \omega, \sin \omega \cos \omega, \sin^2 \omega, \dots$, in the equation obtained in step 2 to zero, we obtain a system of algebraic equations about the unknown numbers $k, \lambda, B_0, A_i, B_i$ for $i = 1, 2, \dots, n$.

Step 4. Using Wu's elimination methods, the algebraic equations in step 3 are solved with the aid of a computer.

These then yield the solitary wave solutions for the system (5).

We remark that the above method yield solutions that include terms $\text{sech} \xi$ or $\tanh \xi$, as well as their combinations. They are different from those that are obtained by other methods, such as the homogenous balance method [5,6].

2 New Explicit Solutions

We assume formal solutions of the form

$$u(x, t) = \varphi(\xi), \xi = \lambda(x - kt + c), \quad (10)$$

where λ, k are constants to be determined later and c_0 is an arbitrary constants. Substituting (10) into (1), we obtain an ordinary differential

$$k\varphi' - p\varphi\varphi' - q\varphi^2\varphi' - \lambda r\varphi'' + s\lambda^2\varphi''' = 0. \quad (11)$$

According to the algorithm described in the previous section, we suppose that (11) has the following formal solutions

$$\varphi(\xi) = A_0 + A_1 \sin \omega + A_2 \cos \omega, \quad (12)$$

and target equation

$$\frac{d\omega}{d\xi} = \sin \omega. \quad (13)$$

With the aid of Mathematica or Maple, from (12) and (13), we can get

$$\begin{aligned} & k\varphi' - p\varphi\varphi' - q\varphi^2\varphi' - \lambda r\varphi'' + s\lambda^2\varphi''' \\ = & [6A_2s\lambda^2 - q(A_2^3 - 3A_1^2A_2)]\sin^4\omega - (6A_1s\lambda^2 + qA_1^3 - qA_1A_2^2)\sin^3\omega\cos\omega \\ & + (2pA_1A_2 + 4qA_0A_1A_2 + 2\lambda rA_1)\sin^3\omega \\ & + [2r\lambda A_2 - P(A_1^2 - A_2^2) - q(2A_0A_1^2 - 2A_0A_2^2)]\sin^2\omega\cos\omega \\ & + [pA_0A_1 - 4A_2s\lambda^2 - kA_2 + q(A_2A_0^2 + A_2^3 - 2A_1^2A_2)]\sin^2\omega \\ & + (s\lambda^2A_1 + kA_1 - pA_0A_1 - qA_1A_0^2)\cos\omega\sin\omega \\ & + [-pA_1A_2 - r\lambda A_1 - 2qA_0A_1A_2]\sin\omega \\ = & 0. \end{aligned}$$

Setting the coefficients of $\sin^j\omega\cos^i\omega$ for $i = 0, 1$ and $j = 1, 2, 3, 4$ to zero, we have the following set of overdetermined equations in the unknowns $A_0, A_1, A_2, \lambda, k$:

$$\begin{aligned} 6A_2s\lambda^2 - q(A_2^3 - 3A_1^2A_2) &= 0, \\ 6A_1s\lambda^2 + q(A_1^3 - 3A_1^2A_2) &= 0, \\ 2pA_1A_2 + 4qA_0A_1A_2 + 2\lambda rA_1 &= 0 \\ 2r\lambda A_2 - p(A_1^2 - A_2^2) - q(2A_0A_1^2 - 2A_0A_2^2) &= 0 \\ pA_0A_1 - 4A_2s\lambda^2 - kA_2 + q(A_2A_0^2 + A_2^3 - 2A_1^2A_2) &= 0, \\ s\lambda^2A_1 + kA_1 - pA_0A_1 - qA_1A_0^2 &= 0, \\ pA_1A_2 + r\lambda A_1 + 2qA_0A_1A_2 &= 0. \end{aligned}$$

We now solve the above set of equations by using the Wu elimination method [4], and obtain the following solution:

$$\text{Case 1. } A_1 = 0, A_2 = \pm\sqrt{\frac{6s}{q}}\lambda, A_0 = \pm\sqrt{\frac{k-2s\lambda^2}{q}}$$

$$k = 12s\lambda + r^2 \pm \frac{1}{2}rp\sqrt{\frac{6s}{q}} + \frac{sp^2}{4q}, sq > 0, q(k - 2s\lambda^2) > 0.$$

If we now take the target equation as

$$\frac{d\omega}{d\xi} = \cos \omega, \quad (14)$$

then proceeding in similar fashions, we get

Case 2. $A_0 = \pm \frac{r}{\sqrt{6qs}} - \frac{p}{2q}$, $A_1 = \pm \sqrt{\frac{6s}{q}}$, $A_2 = 0$,

$$k = \pm \frac{2pr}{\sqrt{6qs}} + \frac{r^2 + 36s^2}{6s} - 4s\lambda^2, qs > 0.$$

Case 3. $A_0 = \pm \frac{r}{\sqrt{6qs}}$, $A_1 = \pm \sqrt{\frac{3s}{2q}}\lambda$, $A_2 = i\sqrt{\frac{3s}{2q}}\lambda$, $qs > 0$,

$$k = \frac{r^2}{6s} \pm \frac{pr}{\sqrt{6qs}} - 4s\lambda^2.$$

Next, integrating $d\omega/d\xi = \sin \omega$ and taking the integration constant zero, we obtain

$$\sin \omega = \operatorname{sech} \xi, \quad (15)$$

and

$$\cos \omega = \pm \tanh \xi. \quad (16)$$

Similarly, from (14), we get

$$\cos \omega = -\operatorname{sech} \xi, \quad (17)$$

$$\sin \omega = \pm \tanh \xi. \quad (18)$$

According to (12), (15)-(18) and the solutions in Cases 1-3, we obtain the following solitary wave solutions of equation (1):

(I) $qs > 0$, $q(k - 2s\lambda^2) > 0$,

$$u_1(x, t) = \pm \sqrt{\frac{k - 2s\lambda}{q}} \pm \sqrt{\frac{6s}{q}} \tanh \lambda(x - kt + c_0),$$

where $k = 12s\lambda^2 + r^2 \pm \frac{1}{12}rp\sqrt{\frac{6s}{q}} + \frac{sp^2}{4q}$.

(II) $qs > 0$,

$$u_2(x, t) = \pm \frac{pr}{\sqrt{6qs}} - \frac{p}{2q} \pm \sqrt{\frac{6s}{q}}\lambda \tanh \lambda(x - kt + c_0),$$

where $k = \pm \frac{2pr}{\sqrt{6qs}} + \frac{r^2 + 36s^2}{6s} - 4s\lambda^2$.

(III) $qs > 0$,

$$u_3(x, t) = \pm \frac{r}{\sqrt{6qs}} + \sqrt{\frac{3s}{2q}}\lambda(\pm i \tanh \xi - \operatorname{sech} \xi)$$

where $\xi = \lambda(x - kt + c_0)$, $k = \frac{r^2}{6s} + \frac{pr}{\sqrt{6qs}} - 4s\lambda^2$.

Note that as $|\xi| \rightarrow +\infty$, $u_3(x, t) \rightarrow \pm[\frac{r}{\sqrt{6qs}} + \sqrt{\frac{3s}{2q}}\lambda]$.

These solutions of (1) are solitary wave solutions. They are linear combinations of kink solitary and bell solitary wave solutions. They are not available in Wang [1] nor in Xia [5].

3 Travelling Wave Solutions

By means of the same procedures described above, we may obtain solutions of (2), (3) and (4):

1. For the compound KdV equation (2), we have the following formal solitary wave solutions.

$$u_4(x, t) = -\frac{p}{2q} \pm \sqrt{\frac{6s}{q}} \tanh \lambda [x - (2s\lambda^2 + \frac{p^2}{4q})t + c_0], \quad qs > 0,$$

$$u_5(x, t) = -\frac{p}{2q} \pm \sqrt{\frac{3s}{2q}} \lambda [\operatorname{sech} \xi \pm i \tanh \xi],$$

where $\xi = \lambda[x - (\frac{1}{2}s\lambda^2 - \frac{p^2}{4q})t + c_0]$ and $qs > 0$.

2. For the KdV-Burgers equation (3), we have the following formal solitary wave solutions

$$u_6(x, t) = \pm \frac{r}{\sqrt{6qs}} \pm \sqrt{\frac{6s}{q}} \lambda \tanh[\lambda(x - \frac{12s^2\lambda^2 + r^2}{6s}t + c_0)], \quad qs > 0,$$

$$u_7(x, t) = \pm [\frac{r}{\sqrt{6qs}} \pm \sqrt{\frac{3s}{2q}} \lambda [\operatorname{sech} \xi \pm \tanh \xi]], \quad qs > 0,$$

where $\xi = \lambda(x - kt + c_0) = \lambda[x - \frac{r^2 + 3s^2\lambda^2}{6s}t + c_0]$.

3. For the mKdV equation (4), we have the the following formal solitary solutions

$$u_8(x, t) = \pm \sqrt{\frac{6s}{q}} \lambda \tanh[\lambda(x - 2s\lambda^2t + c_0)], \quad qs > 0.$$

$$u_9(x, t) = \pm \sqrt{\frac{3s}{2q}} [\tanh \lambda(x - s\lambda^2t + c_0) - i \operatorname{sech} \lambda(x - s\lambda^2t + c_0)], \quad qs > 0.$$

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